

Defects in metals and alloys at finite temperatures from DFT

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ERC CoG project
Materials 4.0

DFG Deutsche
Forschungsgemeinschaft
German Research Foundation

- defects are always there
- defects modify material properties
- need to control defects
- need to **understand** defects



our most favorite *ab initio* simulations

(mine: periodic, plane-wave density functional theory, mostly VASP)

Defects:

point defects (vacancies)

stacking fault energies

anti-phase boundaries

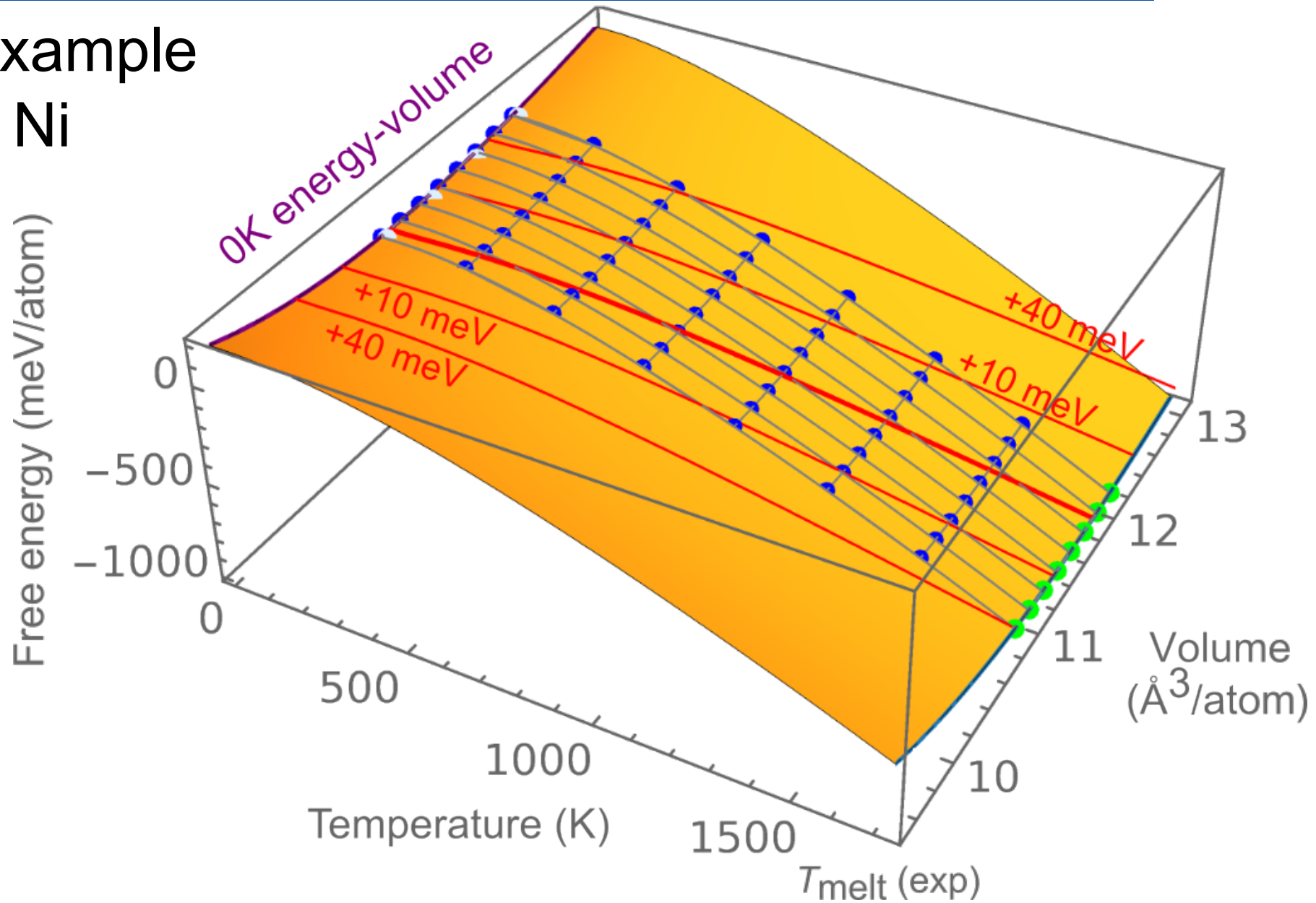
surfaces

- typical experiments at **finite** temperatures
- usual applications at **finite** temperatures
- processing often at **high** temperatures (specifically metals)
- need extension from 0K to **finite** T DFT

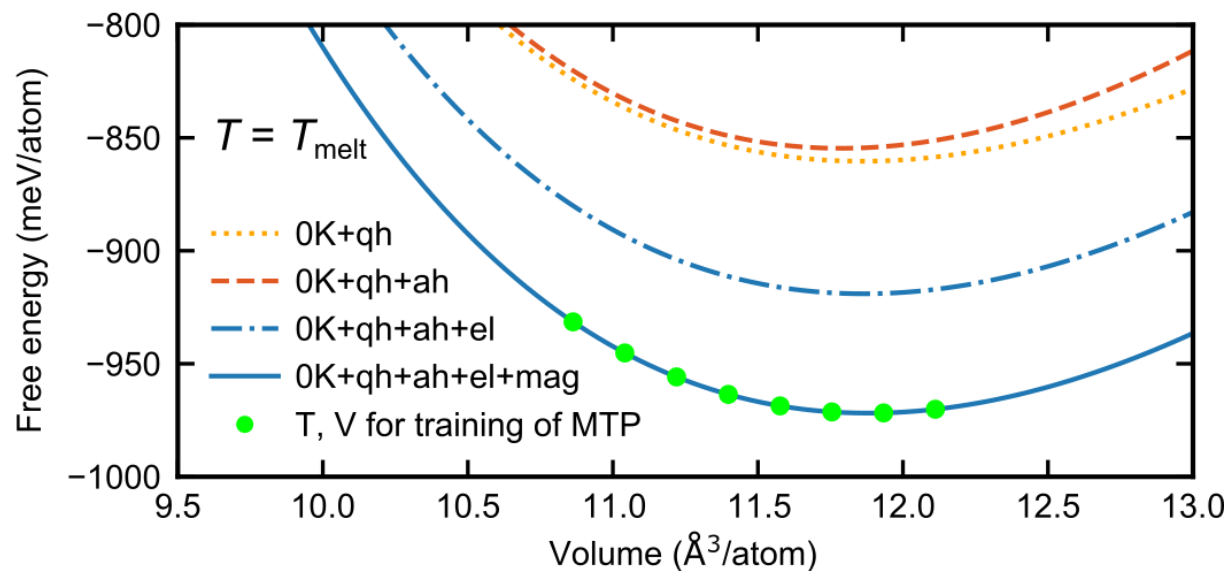
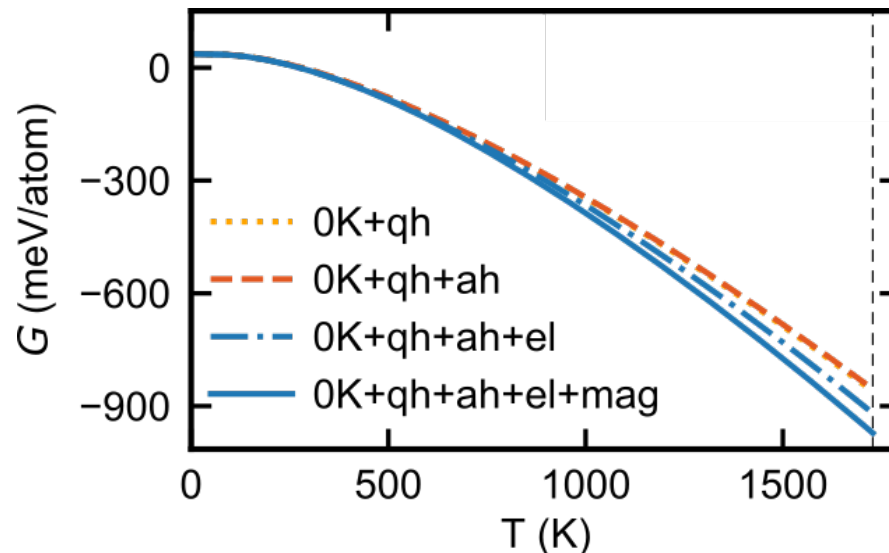
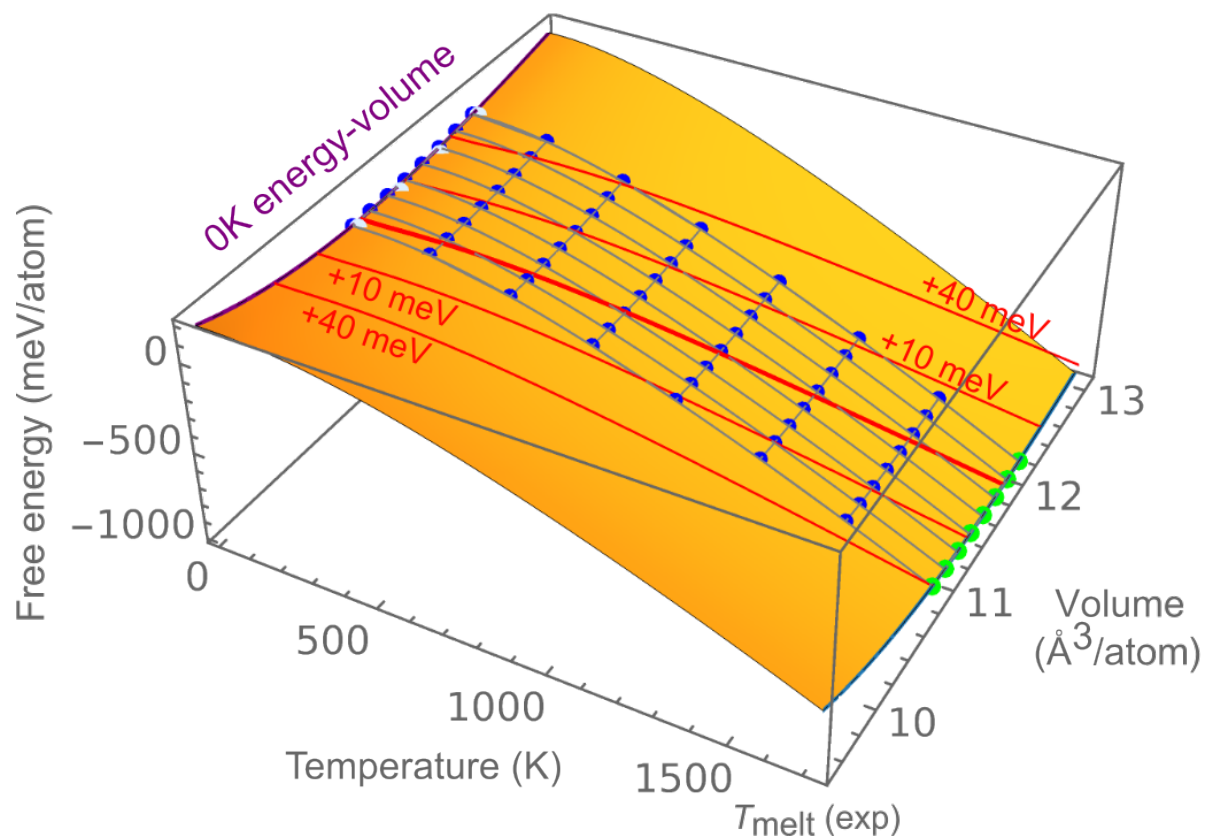
energy → **Helmholtz (free) energy**
Gibbs energy (pressure included)

The free energy surface

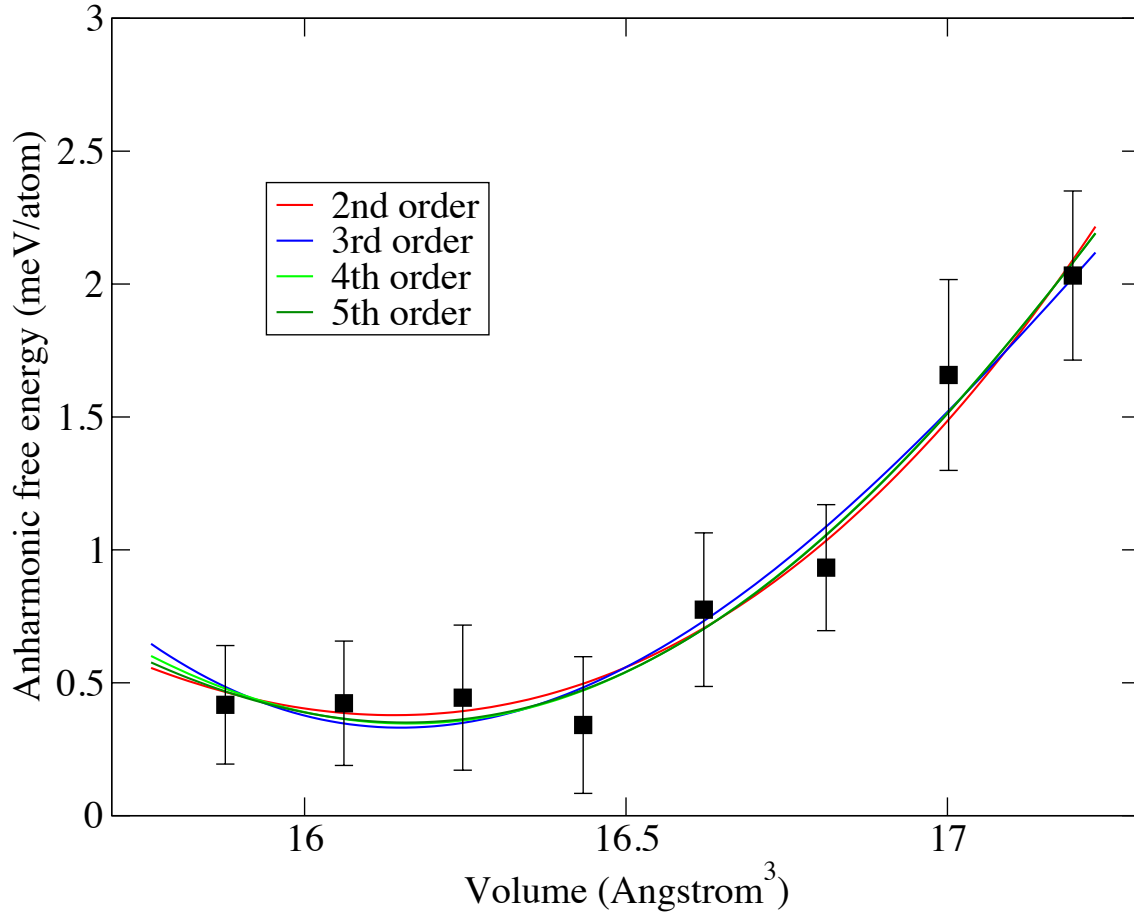
Real-life example for FM fcc Ni



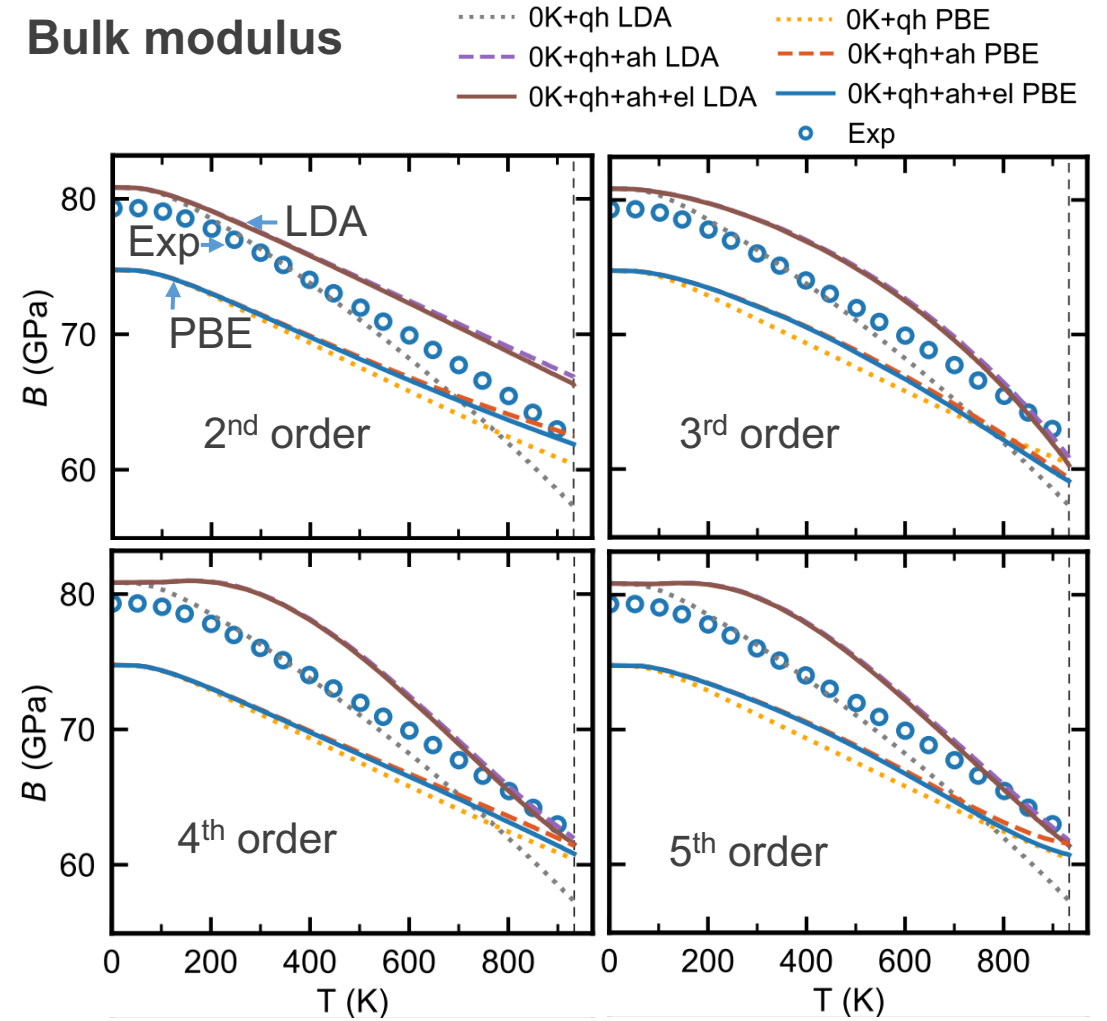
- largest contribution from (quasi)harmonic vibrations
- **BUT: other contributions important for phase stability, thermodynamic properties**



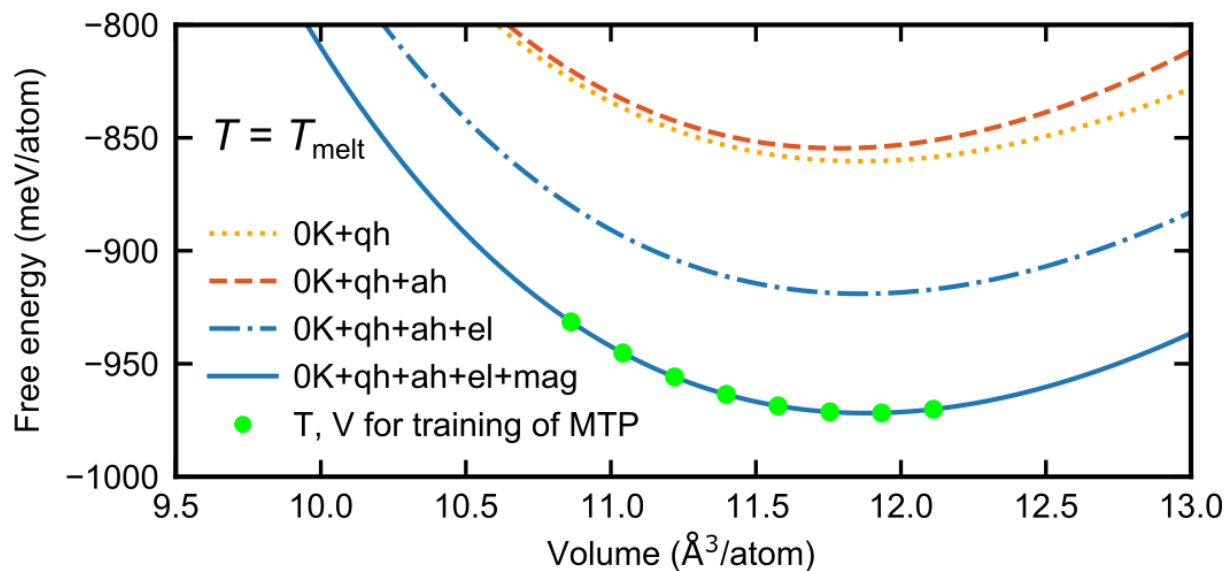
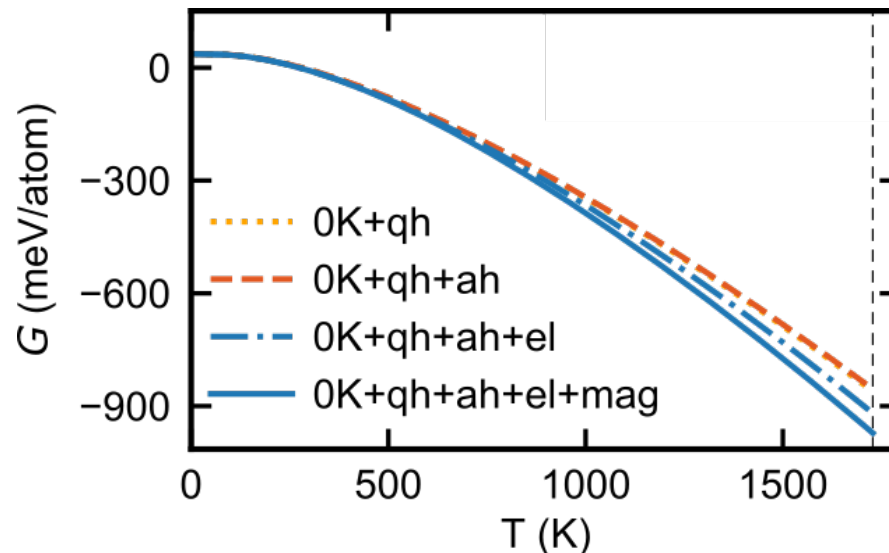
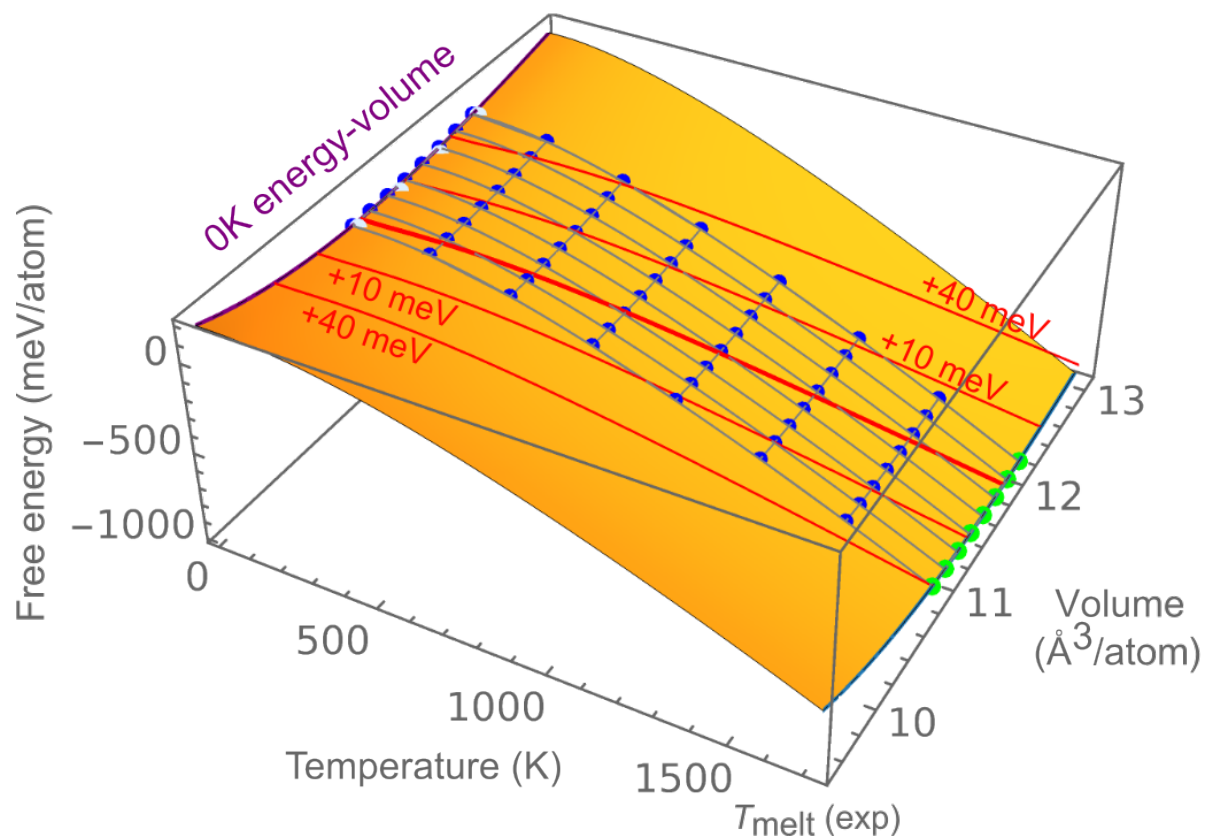
Anharmonic free energy at 933 K for LDA



Bulk modulus

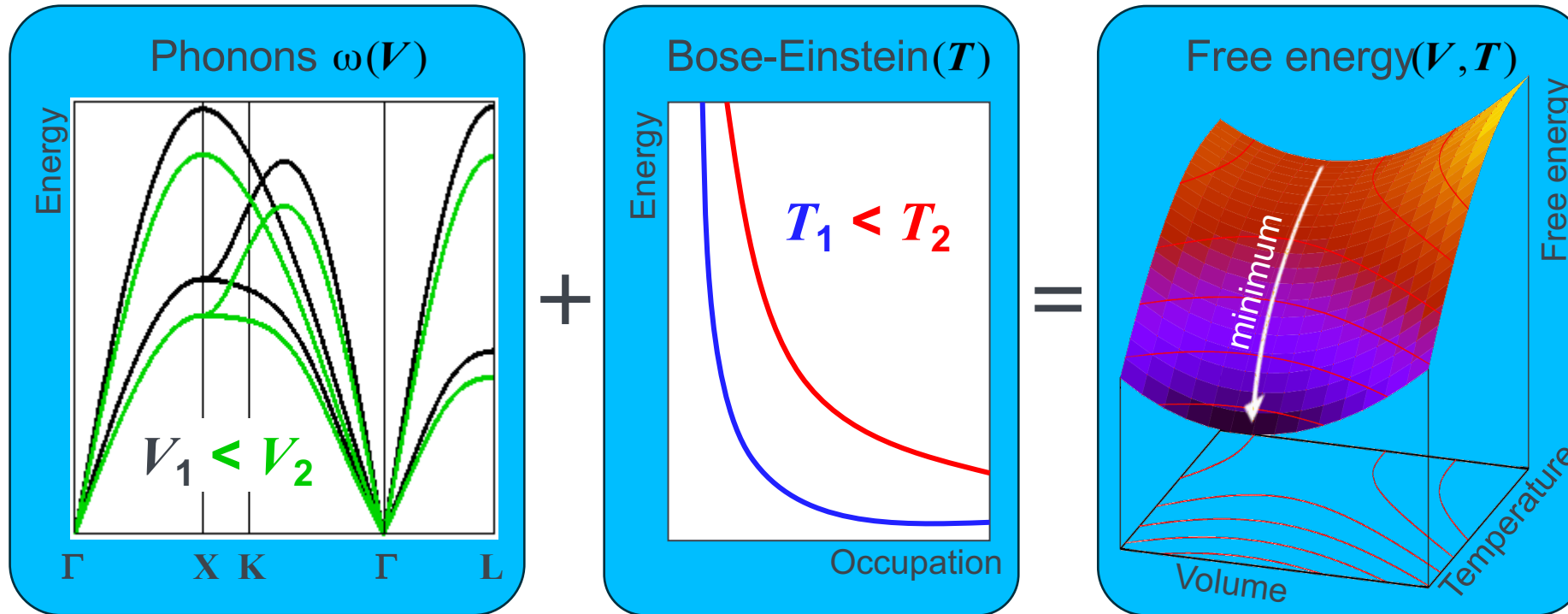


- largest contribution from (quasi)harmonic vibrations
- **BUT: other contributions important for phase stability, thermodynamic properties**



Quasiharmonic vibrations

e.g.: Wallace, *Thermodynamics of Crystals* (Dover, 1998).

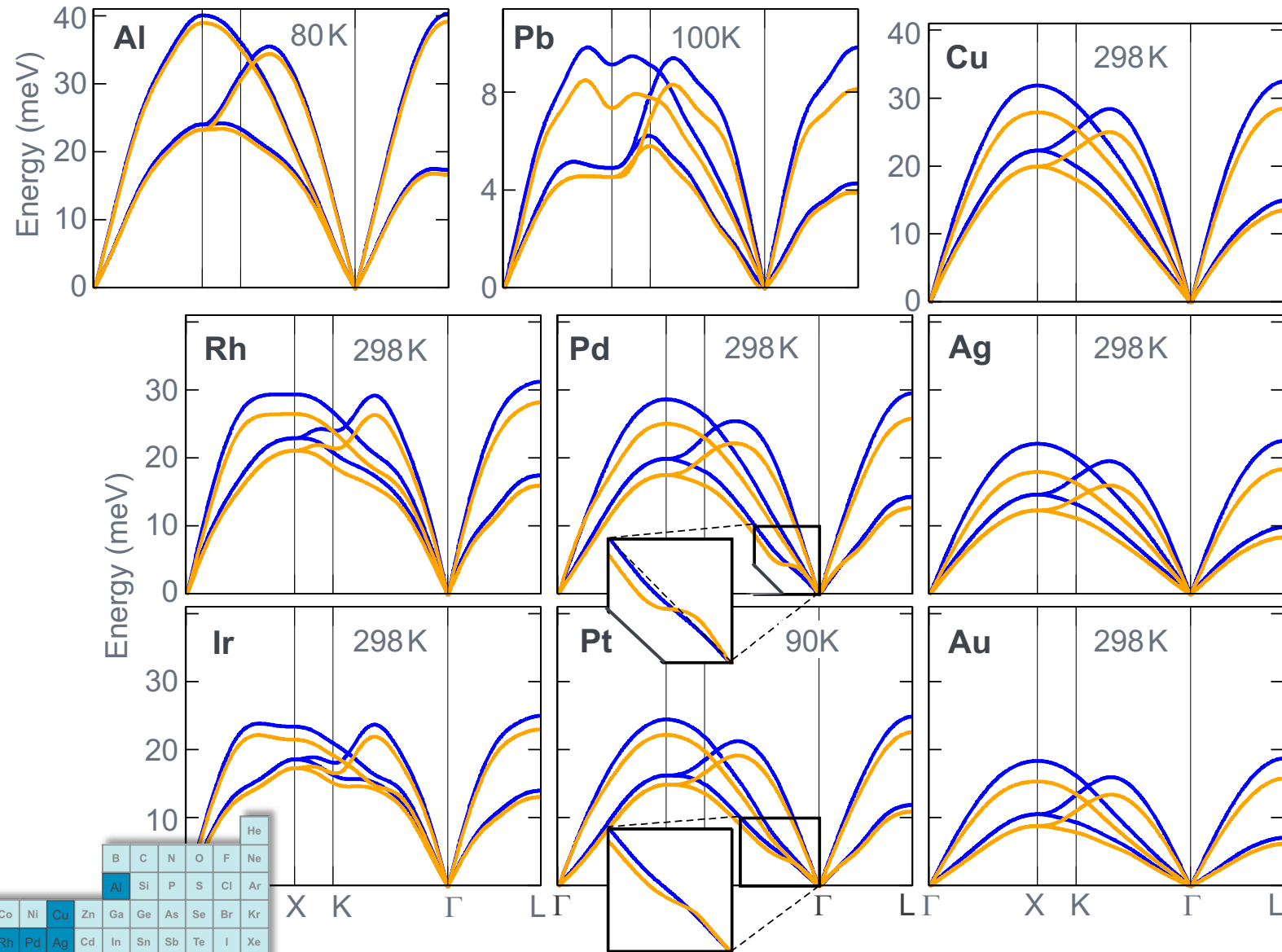
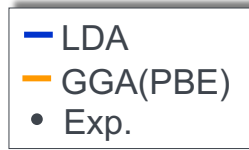
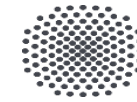


Combination (DFT + quasiharmonic approximation) since ≈ 1989 :

Ref.: Biernacki und Scheffler, *PRL* **63**, 290 (1989).

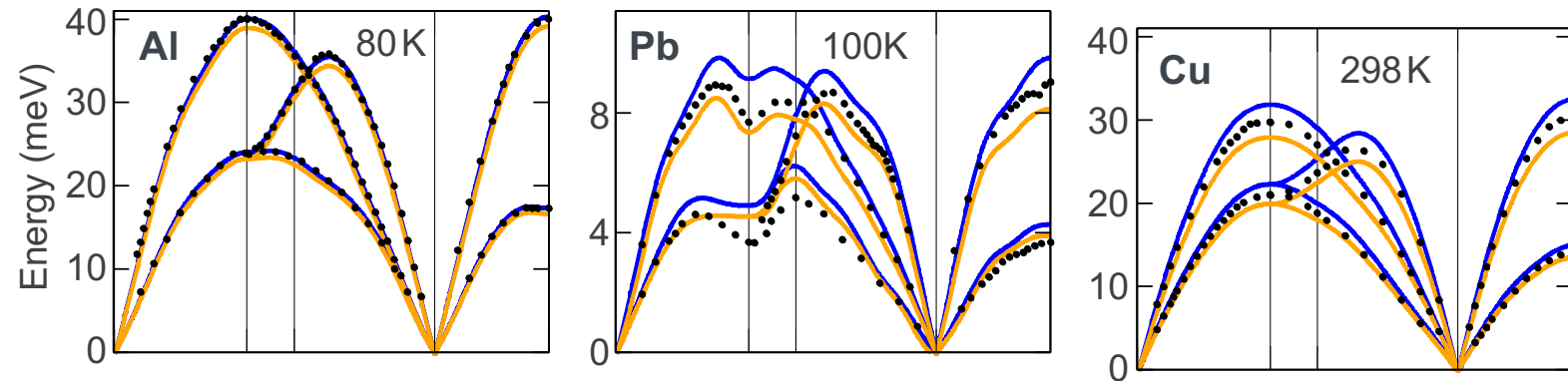
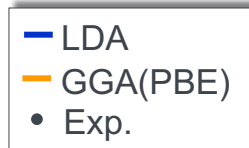
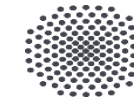
$$F^{\text{qh}}(T, V) = E_{T=0\text{K}}(V) + \sum_{\mathbf{q}} \sum_i^3 \left[\frac{1}{2} \hbar \omega_{\mathbf{q},i}(V) + k_{\text{B}} T \ln \left[1 - \exp \left(-\frac{\hbar \omega_{\mathbf{q},i}(V)}{k_{\text{B}} T} \right) \right] \right]$$

Phonons for fcc metals

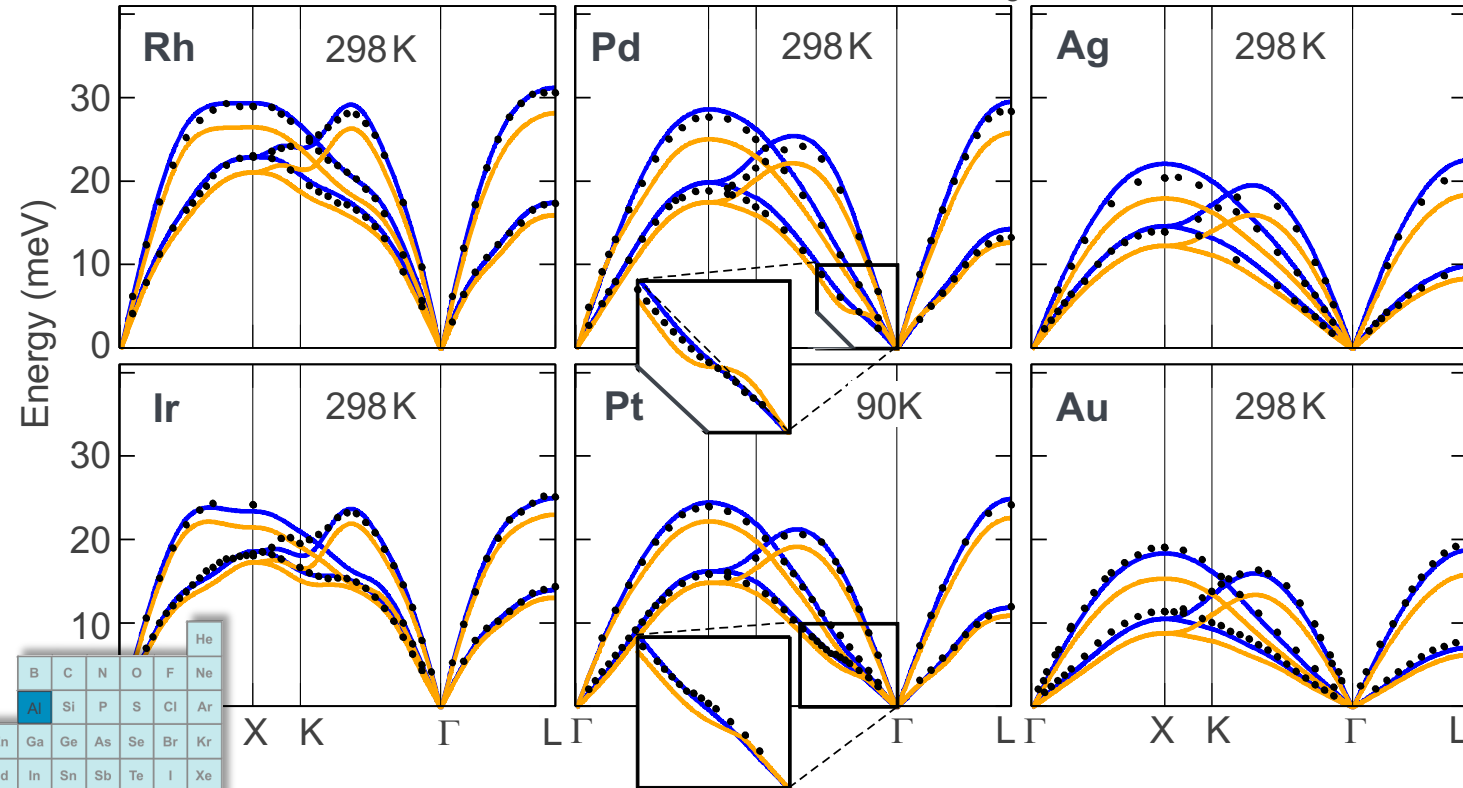


H																	He				
Li	Be															B	C	N	O	F	Ne
Na	Mg											Al	Si	P	S	Cl	Ar				
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr				
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe				
Cs	Ba	Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn				
Fr	Ra	Lr	Rf	Db	Sg	Bh	Hs	Mt	Ds	Uuu	Uub										

Phonons for fcc metals



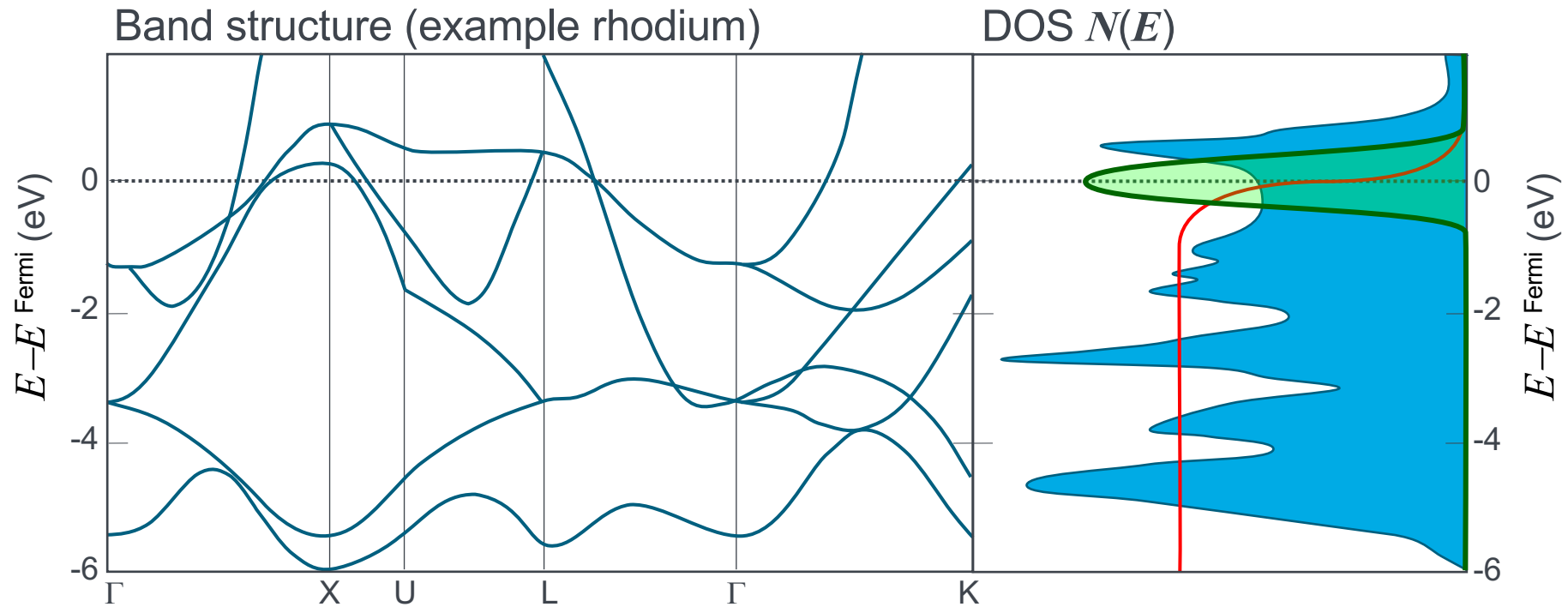
→ *Ab initio*
confidence
interval!



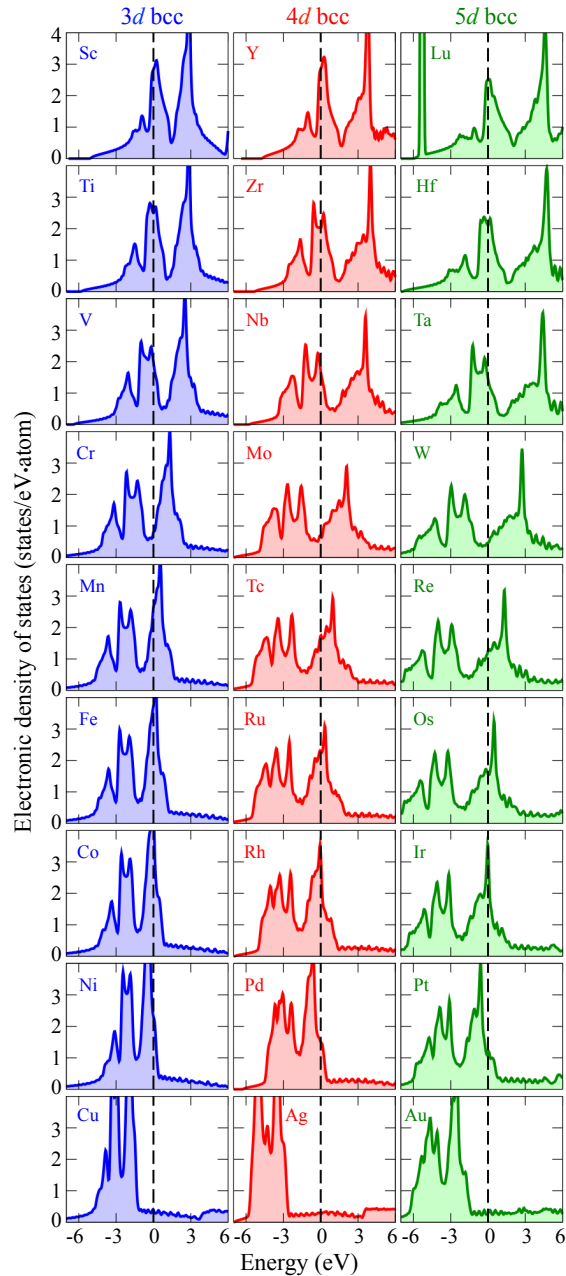
H																	He				
Li	Be															B	C	N	O	F	Ne
Na	Mg															Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr				
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe				
Cs	Ba	Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn				
Fr	Ra	Lr	Rf	Db	Sg	Bh	Hs	Mt	Ds	Uuu	Uub										

Electronic excitations

Mermin, Phys. Rev. 137, A1441 (1965)



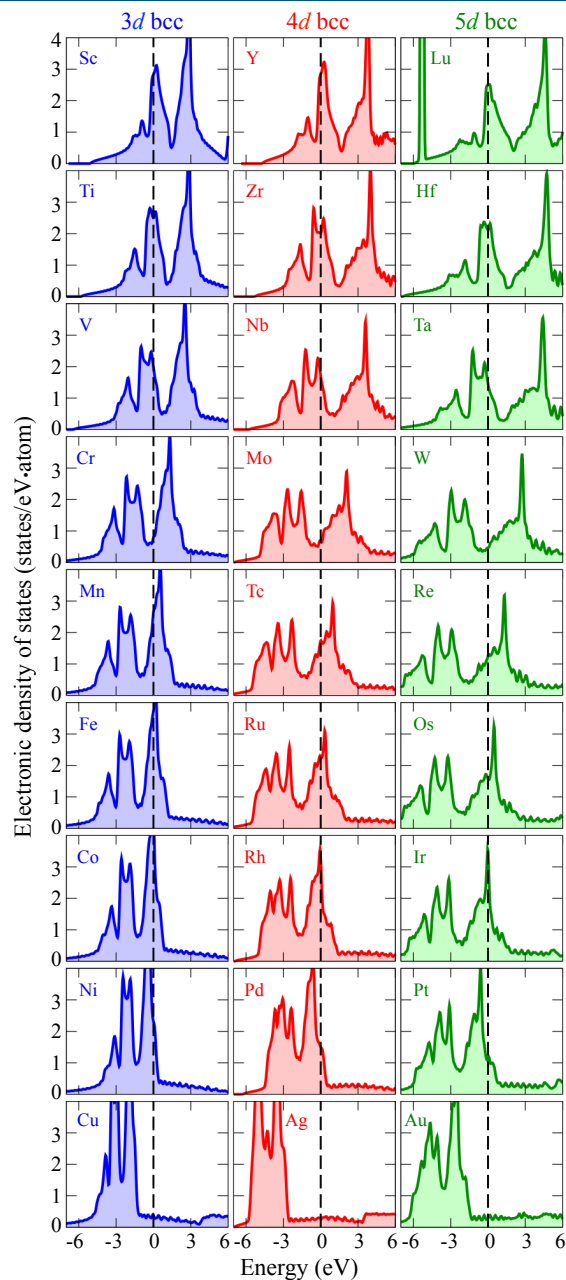
$$F^{\text{el}} = k_{\text{B}}T \int \overset{\text{DOS}}{\downarrow} N(E) \underbrace{\left[\overset{\text{Fermi-Dirac}}{\downarrow} f \ln f + (1 - f) \ln(1 - f) \right]}_{\text{ideal entropy}} dE$$



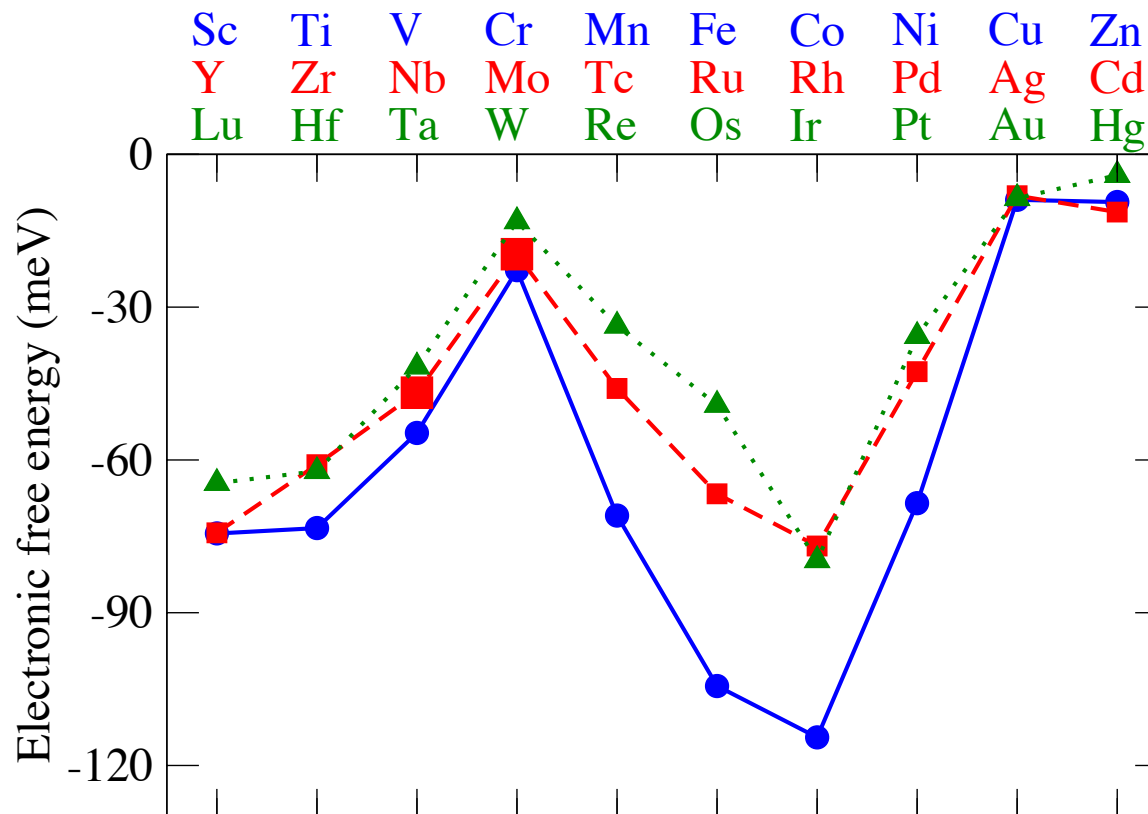
Trends in the periodic table

1A												8A						
1 H 1.00794	2A											5		6	7	8	9	10
3 Li	4 Be											5 B	6 C	7 N 14.0067	8 O 15.9994	9 F 18.9984032	10 Ne 20.1797	
11 Na	12 Mg											13 Al	14 Si	15 P 30.973762	16 S 32.065	17 Cl 35.453	18 Ar 39.948	
19 K	20 Ca	3B	4B	5B	6B	7B	8B		1B	2B	31 Ga	32 Ge	33 As 74.92160	34 Se 78.96	35 Br 79.904	36 Kr 83.798		
37 Rb 85.4678	38 Sr 87.62	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In 114.818	50 Sn 118.710	51 Sb 121.760	52 Te 127.60	53 I 126.90447	54 Xe 131.293	
55 Cs 132.9054519	56 Ba 137.327	57-71 Lanthanides	72 Lu	73 Hf	74 Ta	75 W	76 Re	77 Os	78 Ir	79 Pt	80 Au	81 Tl 204.3833	82 Pb 207.2	83 Bi 208.98040	84 Po [209]	85 At [210]	86 Rn [222]	
87 Fr [223]	88 Ra [226]	89-103 Actinides	104 Rf [267]	105 Db [268]	106 Sg [271]	107 Bh [272]	108 Hs [270]	109 Mt [276]	110 Ds [281]	111 Rg [280]	112 Cn [285]	113 Uut [284]	114 Fl [289]	115 Uup [288]	116 Lv [293]	117 Uus [294]	118 Uuo [294]	

Phys. Rev. B 95 (2017) 165126



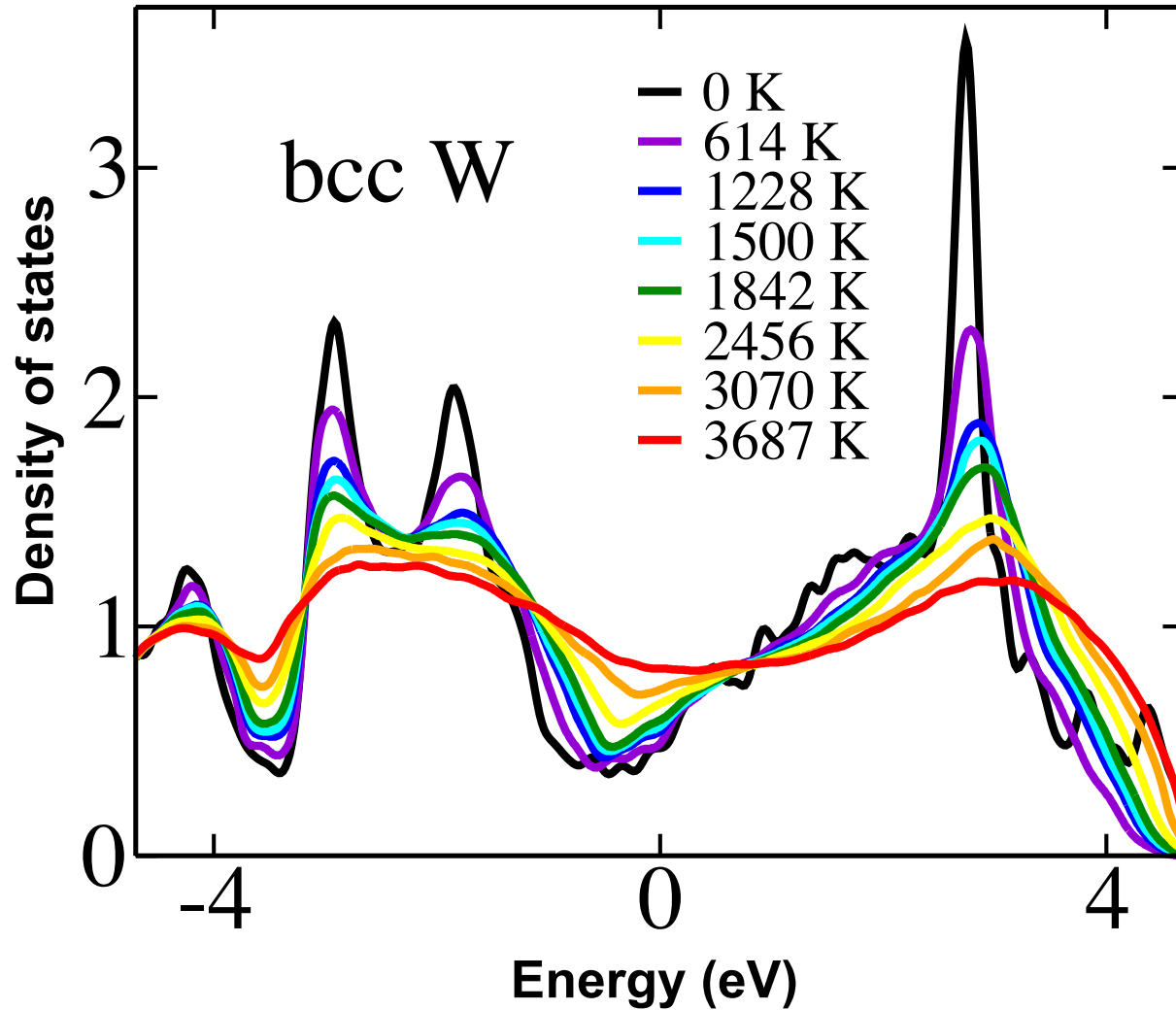
Electronic free energies at 1500 K



→ Electronic free energies
important for phase stabilities

Phys. Rev. B 95 (2017) 165126

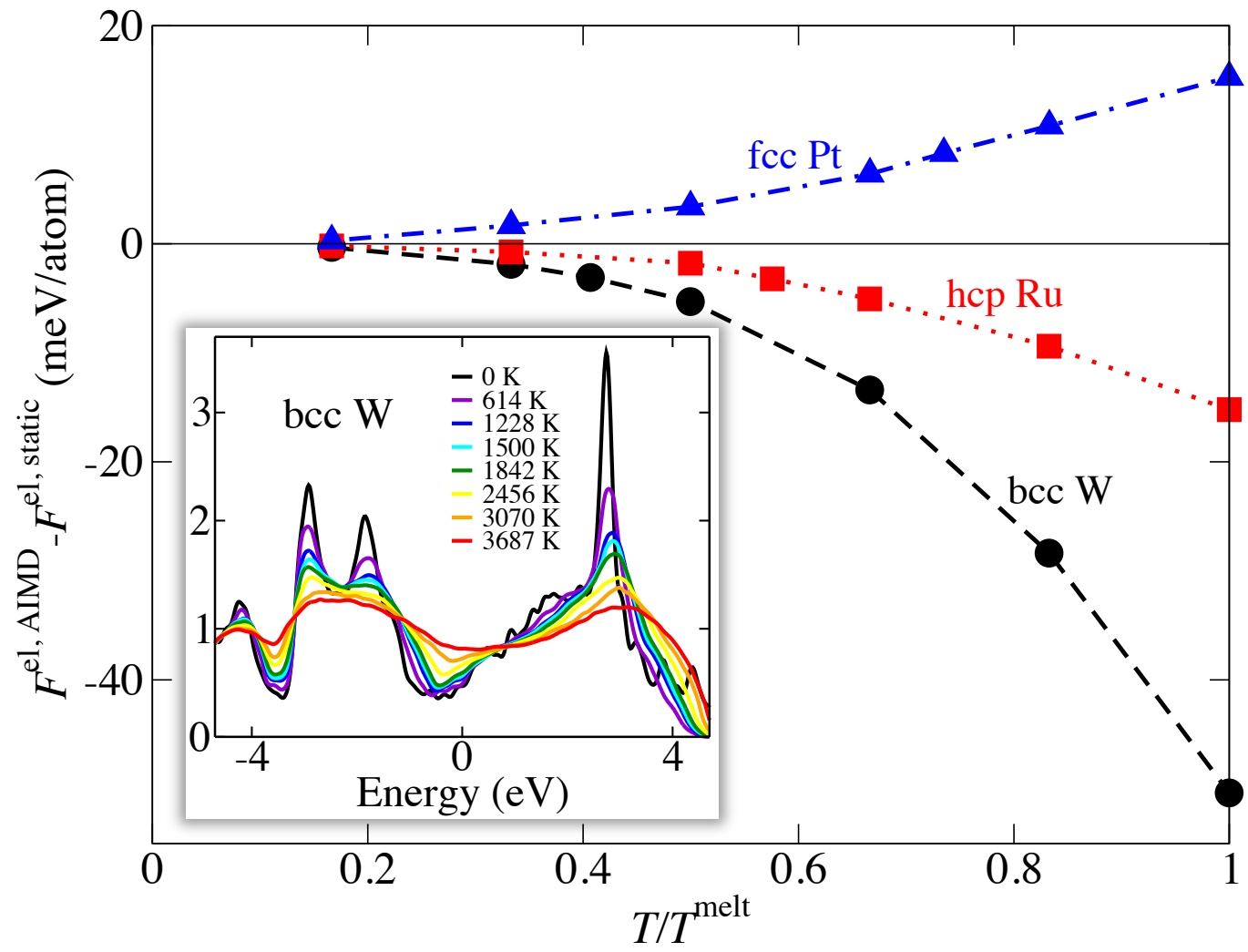
Example for tungsten



- Very strong broadening
- Thermal disorder
- Localization

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Impact on electronic free energies

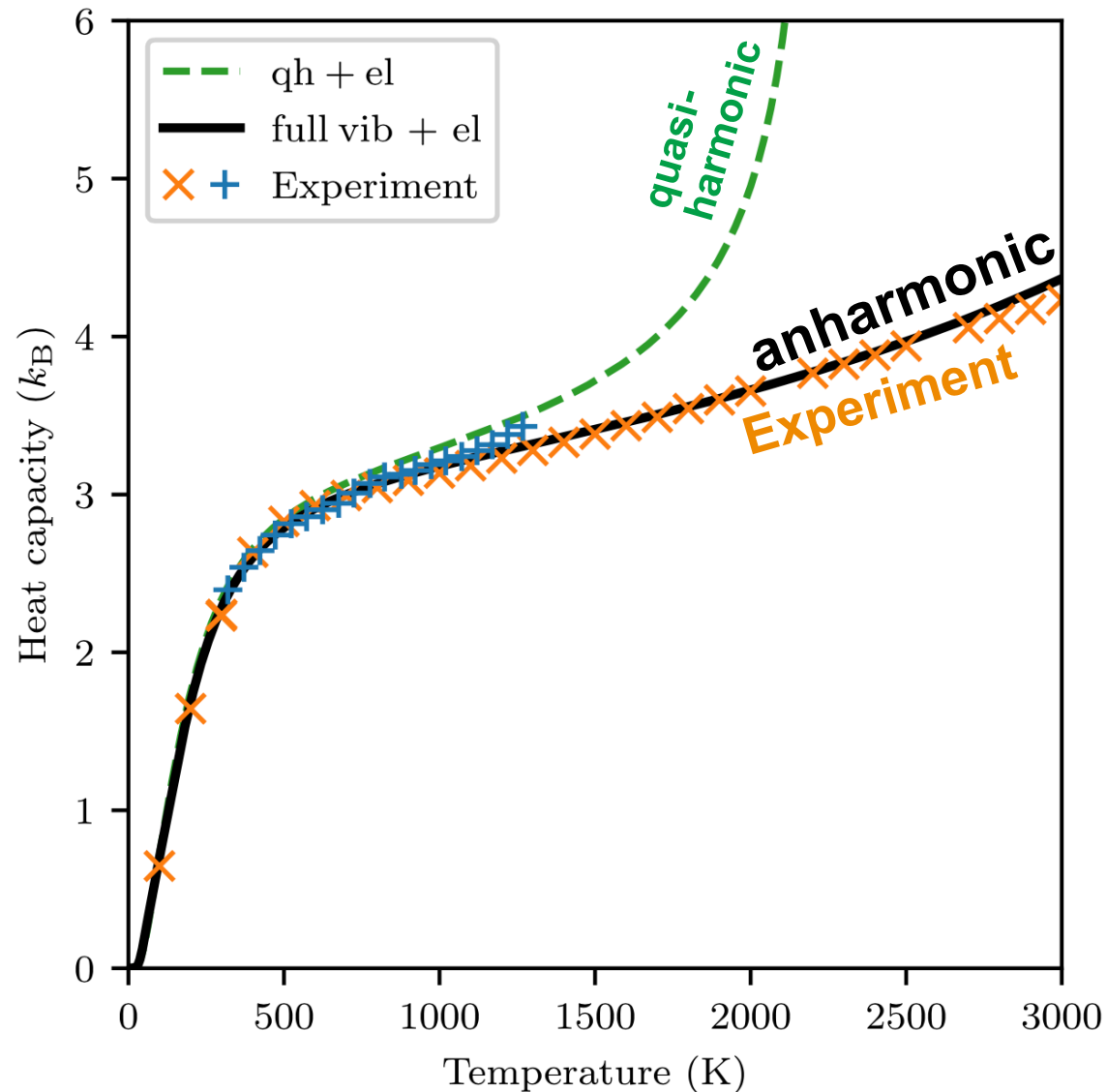


→ Very strong coupling

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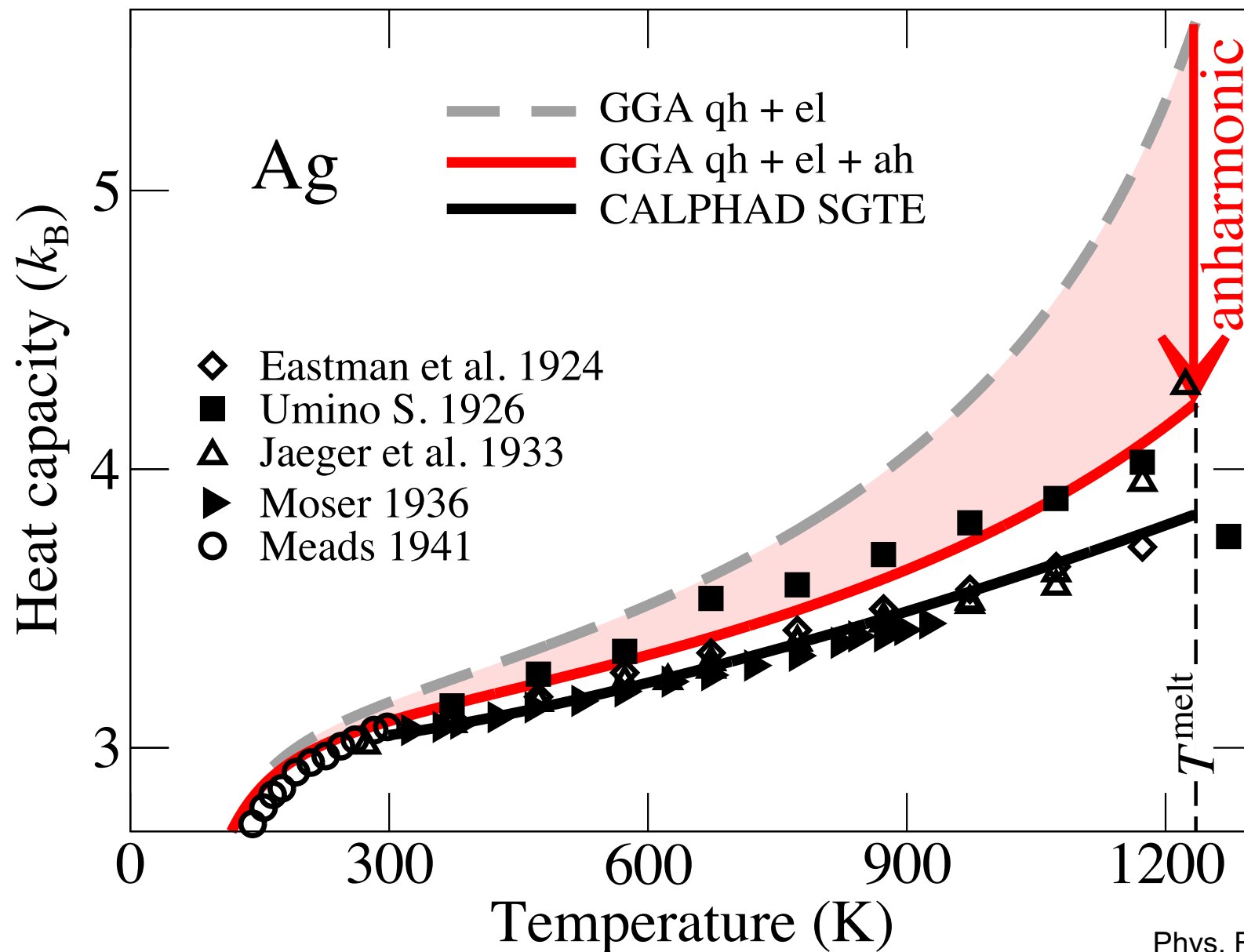
Anharmonic vibrations

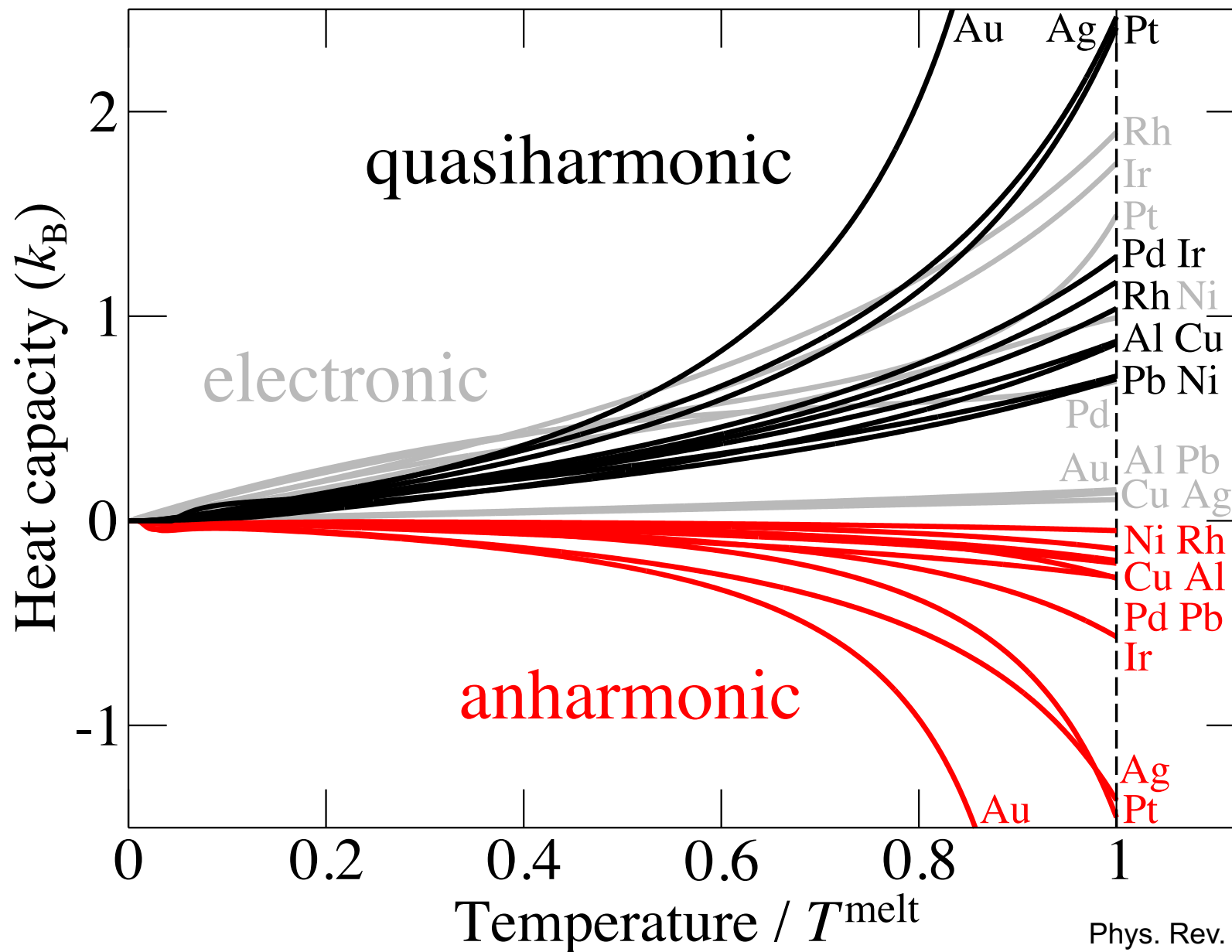
Example: Heat capacity of TiN



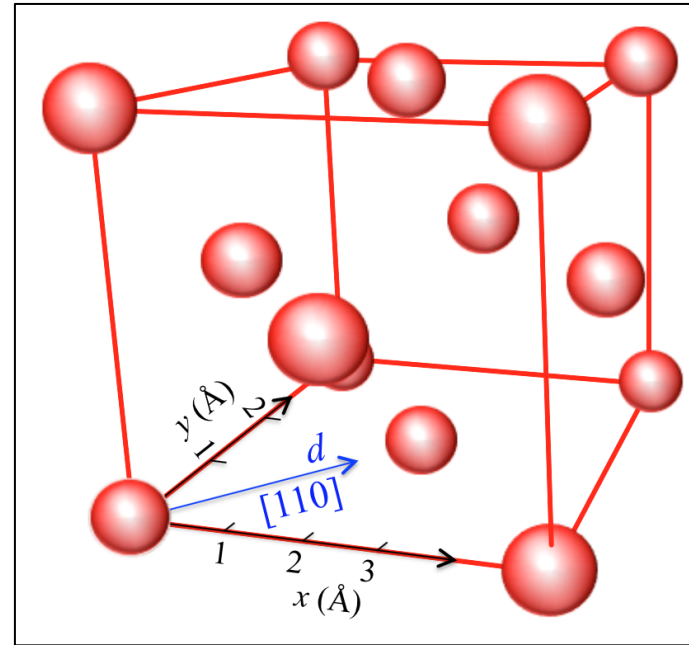
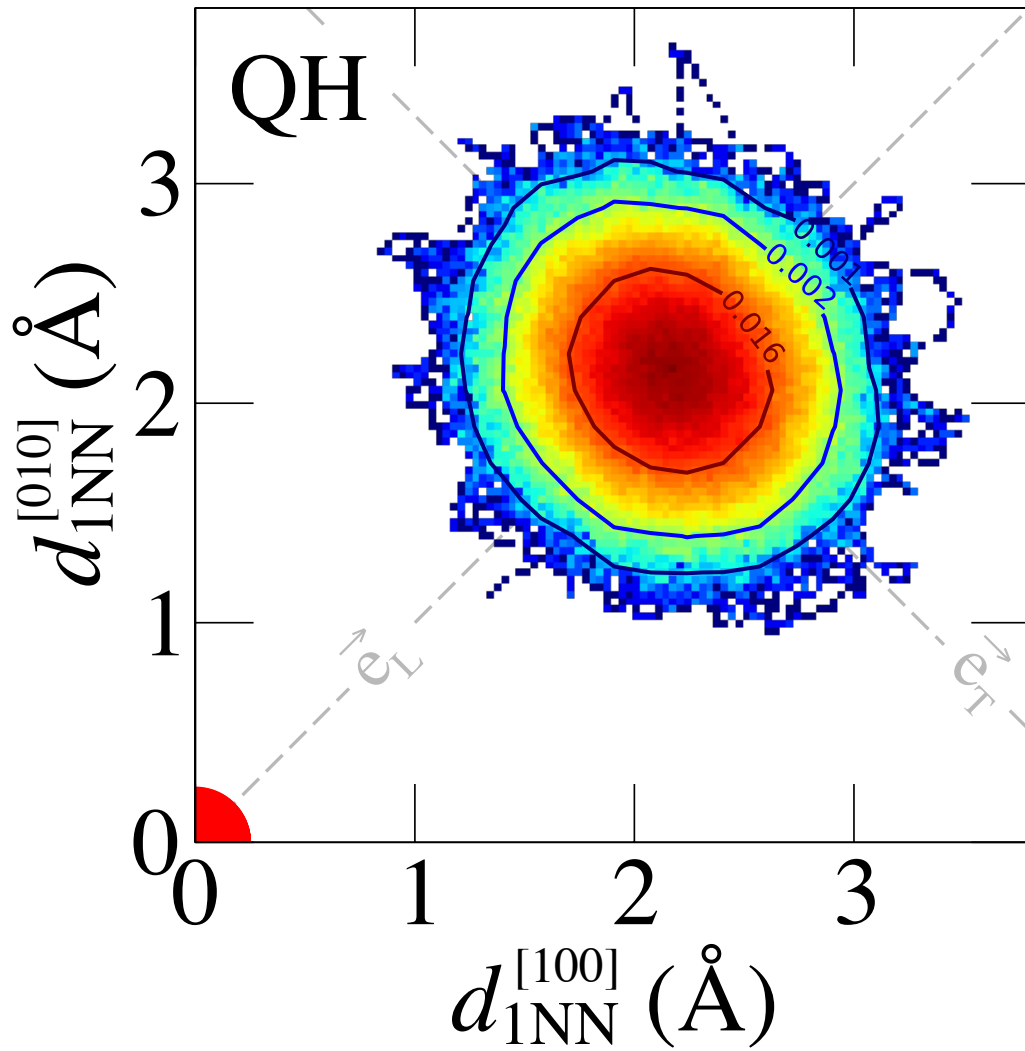
- anharmonic (phonon-phonon interaction) contribution can be very important
- many such examples
- **anharmonicity cannot be neglected in general!**

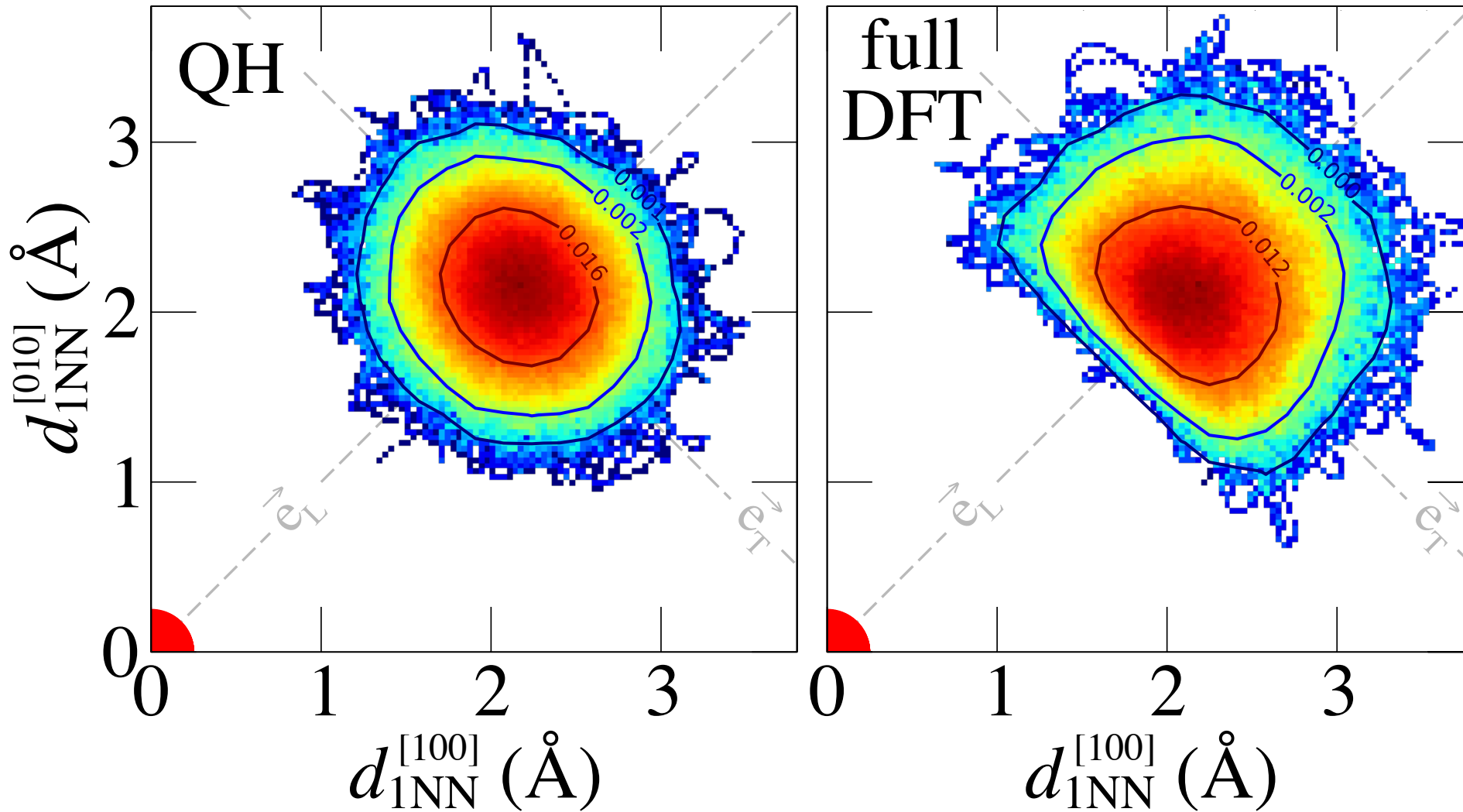
Reference:
Forslund, Grabowski et al.,
PHYSICAL REVIEW B 103,
195428 (2021)



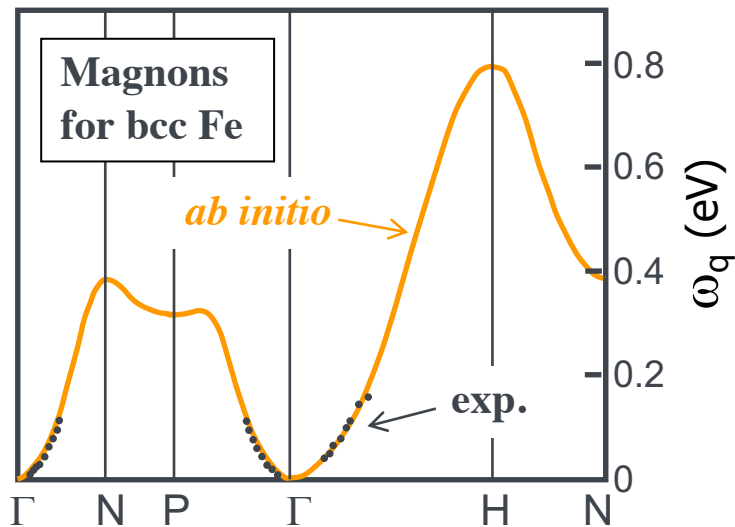
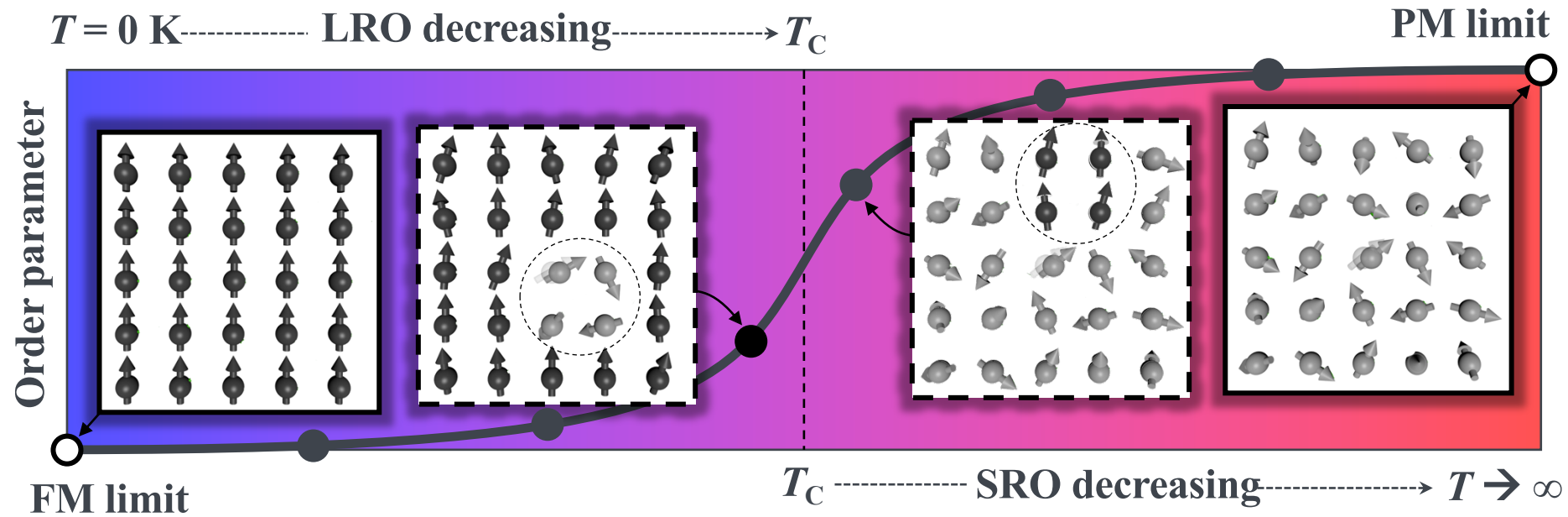


Phys. Rev. Lett. 114 (2015) 195901





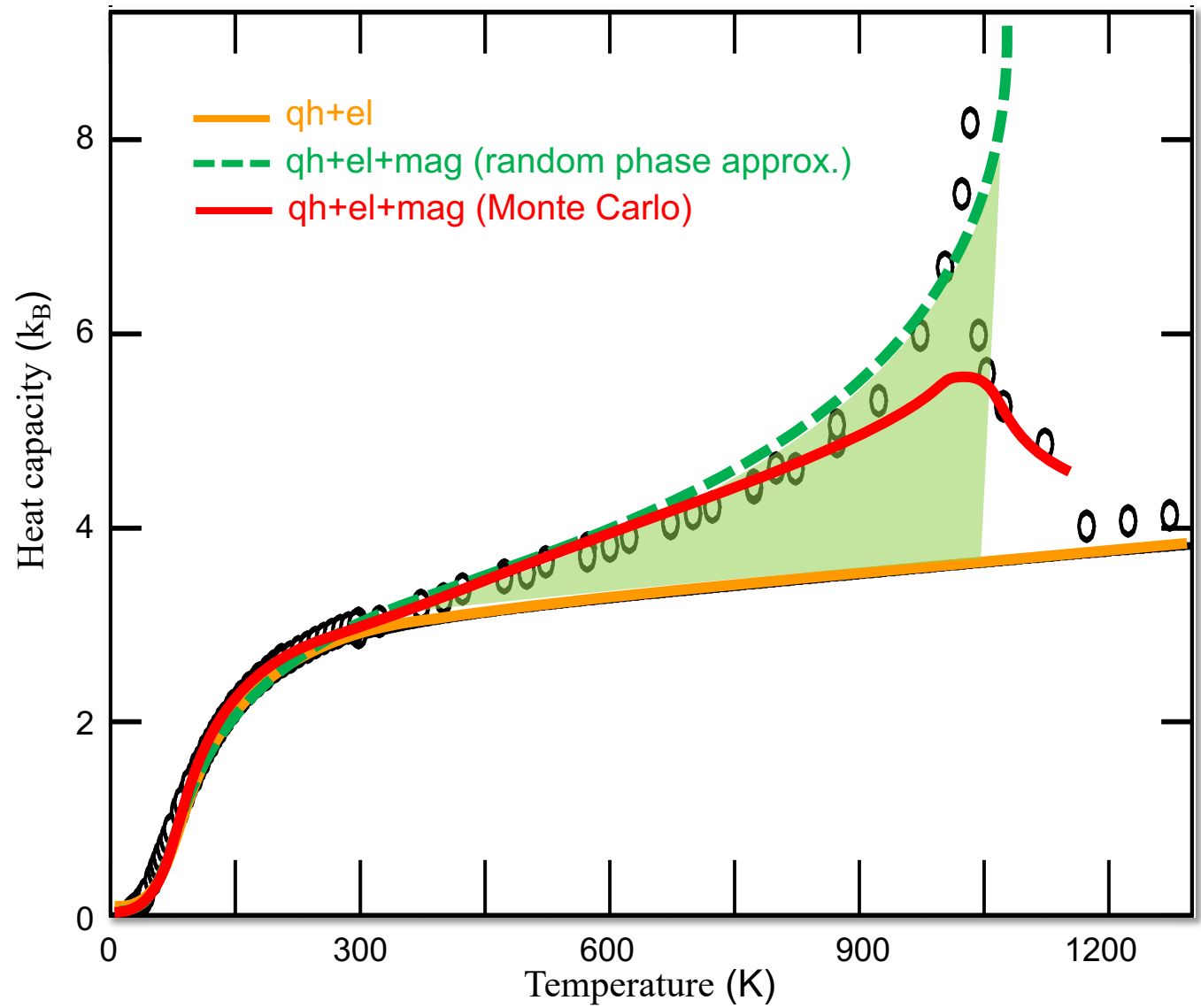
Magnetic excitations



Phys. Rev. B
78 (2008)
033102

- Magnons are bosons
- BUT no simple occupation due to strong interactions
- BETTER mean field theories

Example: Heat capacity for iron

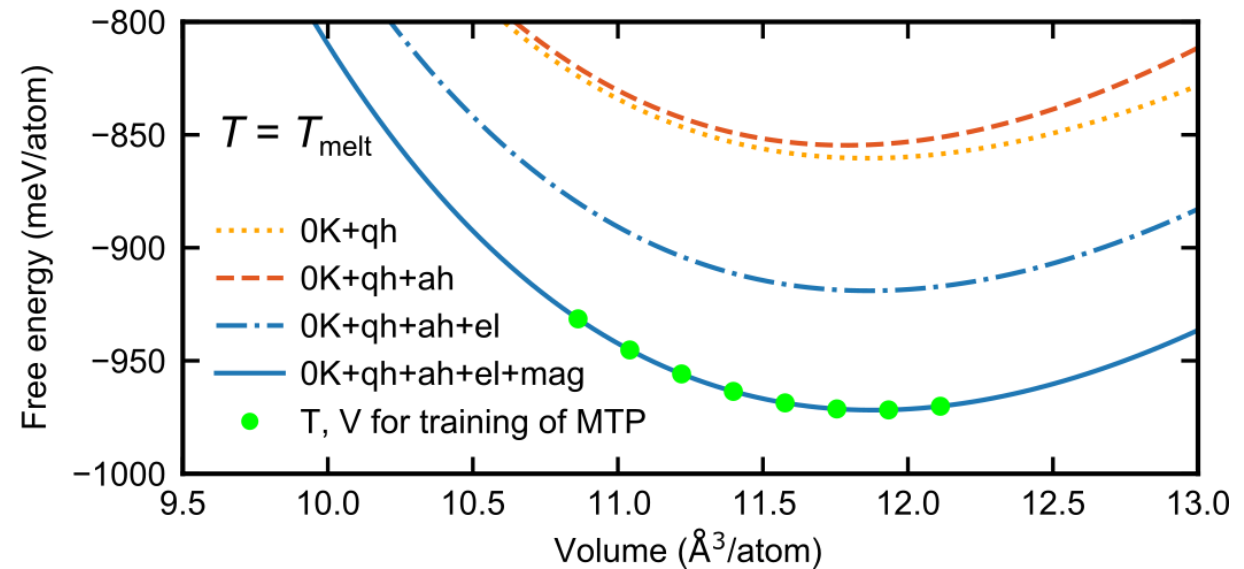
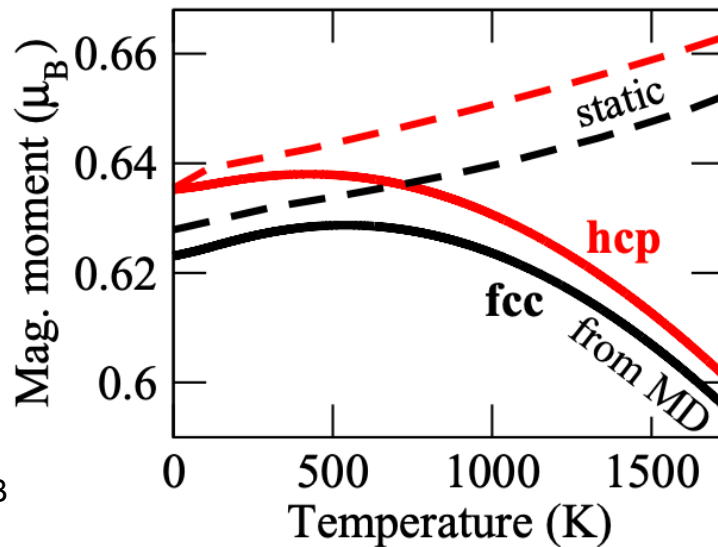


Phys. Rev. B
78 (2008)
033102

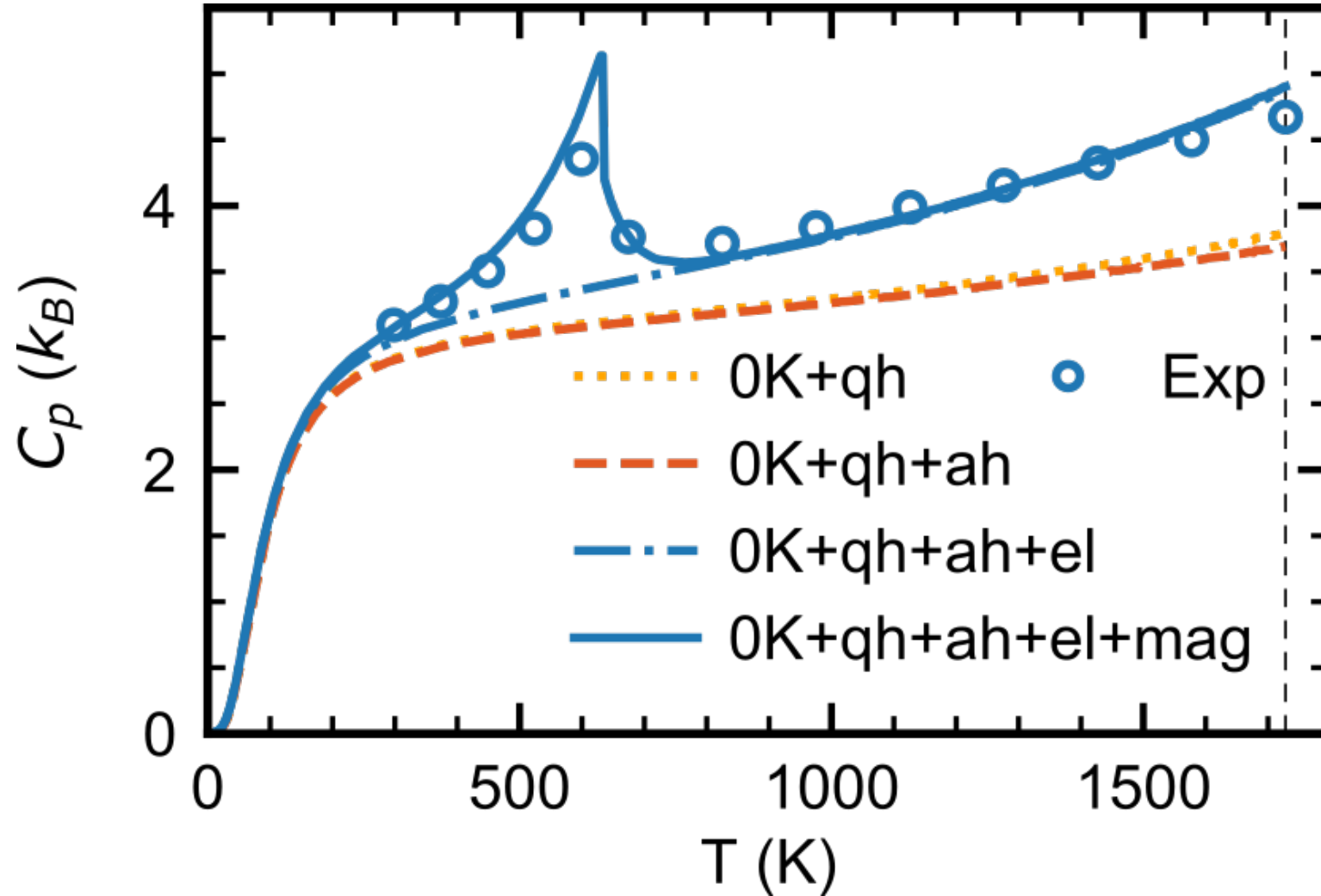
- mean-field approximation if well above Curie temperature
- or empirical approach with DFT/experimental input

$$C^{\text{mag}}(T) = \begin{cases} k_f(T/T_C) \exp[-4(1 - T/T_C)] & T < T_C \\ k_p(T/T_C) \exp[8p(1 - T/T_C)] & T > T_C \end{cases}$$

$$S^{\text{mag}} = k_B \ln(m + 1)$$



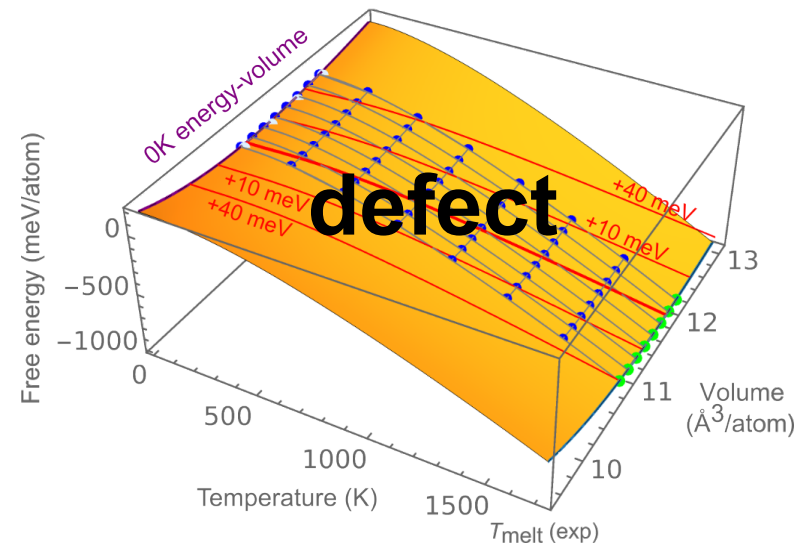
Example: Heat capacity for nickel



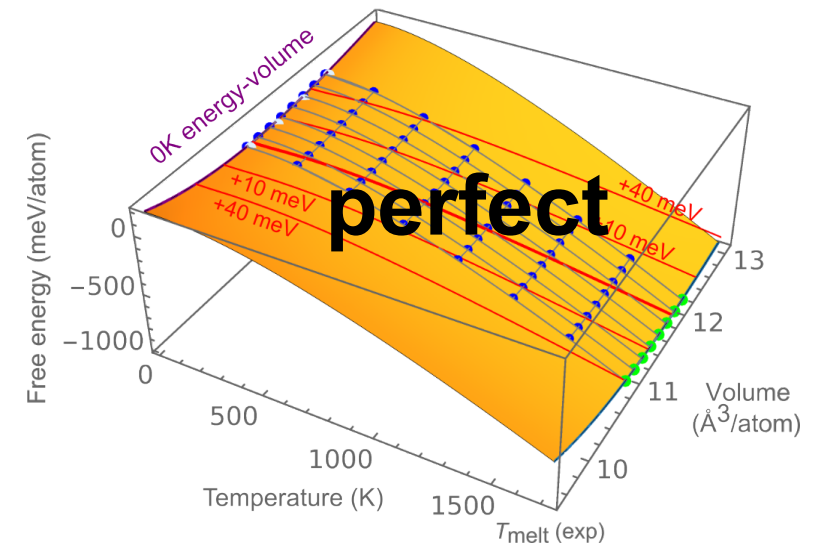
Back to defects

“master” equation:

**defect
formation
free energy** =



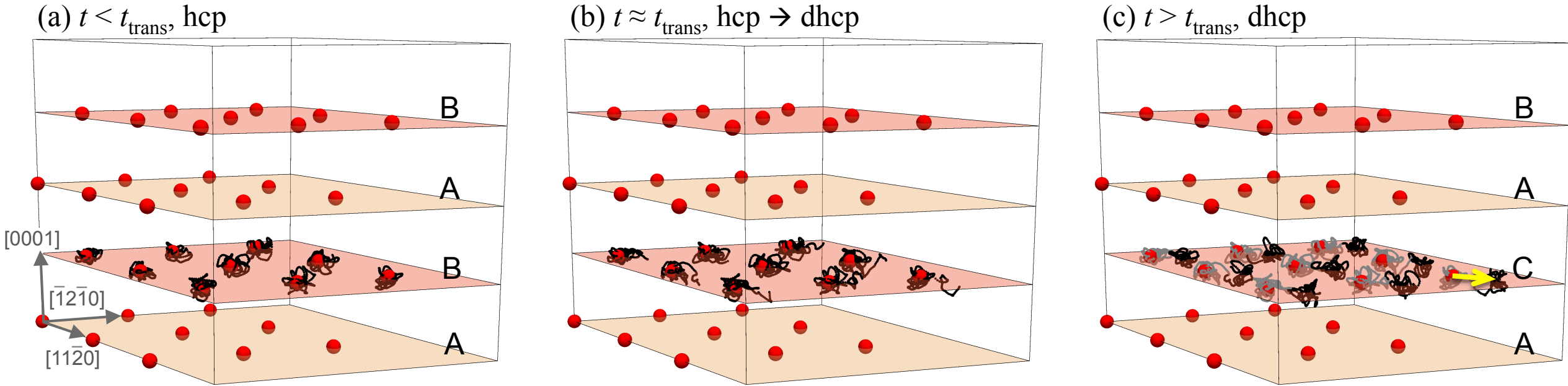
−



- once you have anharmonic vibrations, do not care much about defects
- but supercell size convergence can be tough (error scales with # atoms)
- and sometimes special care necessary, e.g.:
 - stacking plane migration
 - vacancy migration

ANNNI = axial-next-nearest- neighbor-Ising

Example: hcp Ni for ANNNI model ($SFE = hcp - fcc$)



Phys. Rev. B 98 (2018) 224106

- \rightarrow MD needs to be monitored carefully
- \rightarrow here finite size effects play a role
- \rightarrow supercell size needs to be large enough

**Most recent method
for finite temperature DFT**

Methodology: Direct upsampling

(to be published)

1. Run low accuracy DFT molecular dynamics
2. Fit preliminary moment tensor potential (MTP)
3. Optimize MTP on high accuracy DFT snapshots
4. TILD from effective harmonic to MTP for $F(V,T)$
5. Upsampling from MTP to DFT on snapshots for $F(V,T)$
6. Parametrize $F(V,T) \rightarrow$ Legendre transform to $G(P,T)$
7. Extract thermodynamic properties
(e.g. heat capacity, expansion coefficient, bulk modulus)

TILD = thermodynamic integration using Langevin dynamics

Shapeev, Multiscale Modeling & Simulation 14, 1153 (2016)

Descriptors of atomic environments:

- Moments of inertia of surrounding atoms
- They satisfy the needed symmetries (rotation, permutation, translation, ...)

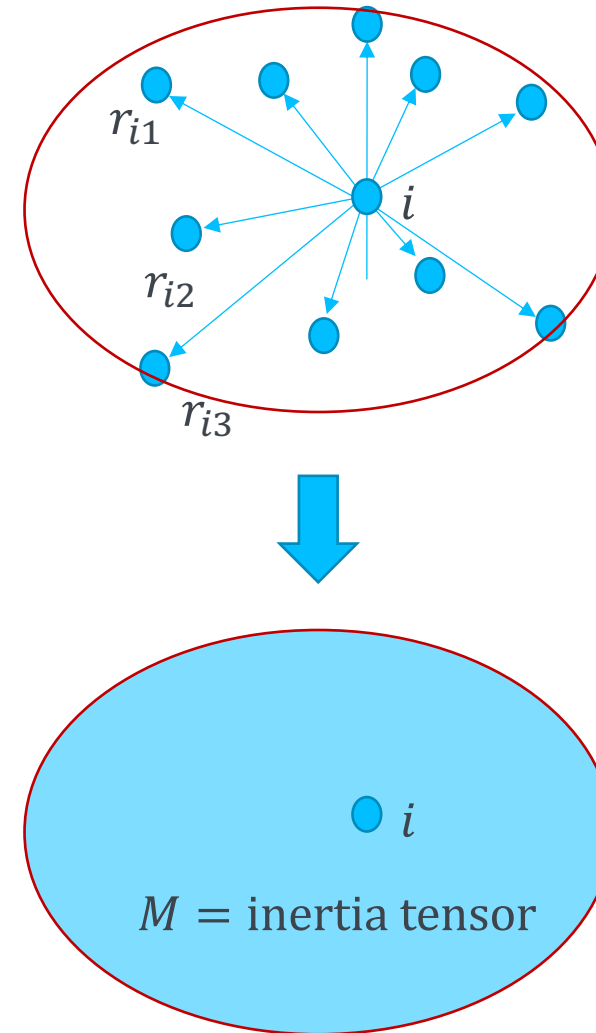
• **Math:**

➤ Moments: $M_{n,m}(\mathbf{r}_i) = \sum_j f_n(|r_{ij}|) \underbrace{r_{ij} \otimes \dots \otimes r_{ij}}_{m \text{ times}}$

➤ Basis functions B_α are (all) different multiplications (contractions) of inertia tensors $M_{m,n}$ yielding a scalar

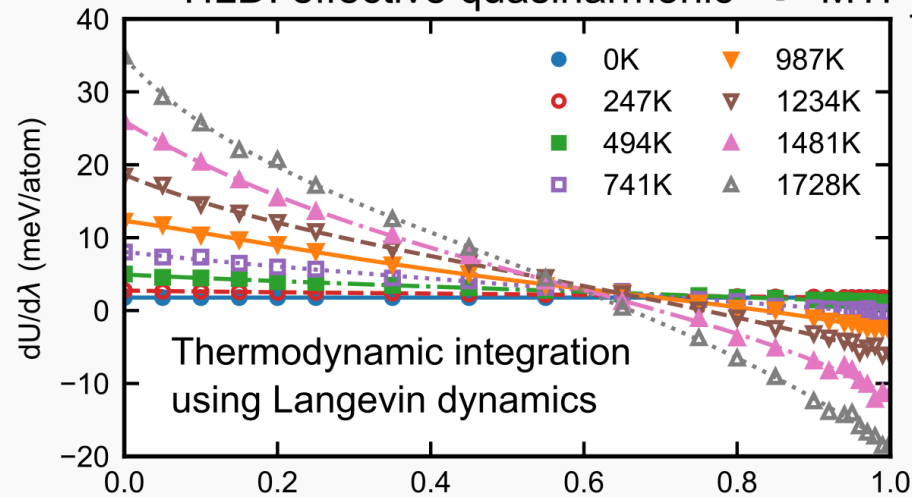
➤ Energy: $E = \sum_\alpha \theta_\alpha B_\alpha$

➤ Active learning via D-optimality

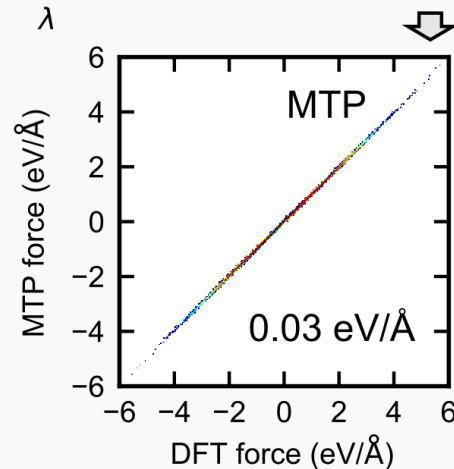
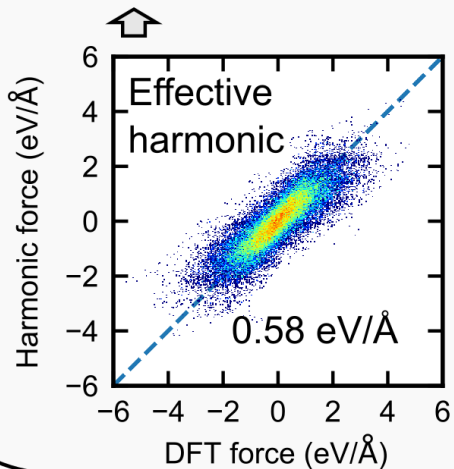


Direct upsampling: Efficient computational method based on machine learning potentials

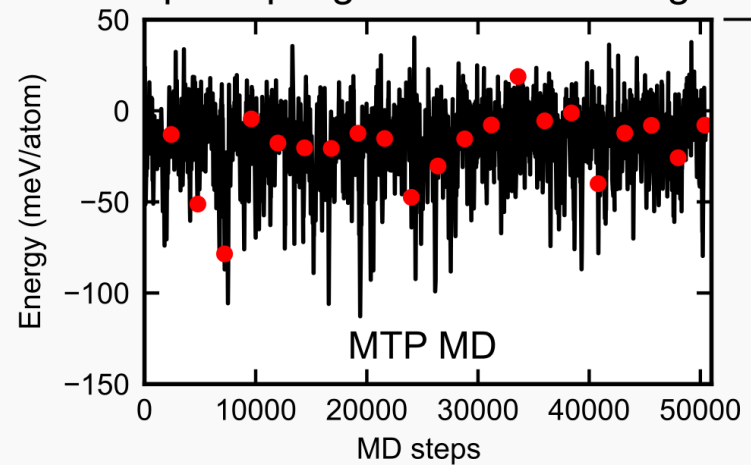
TILD: effective quasiharmonic \rightarrow MTP



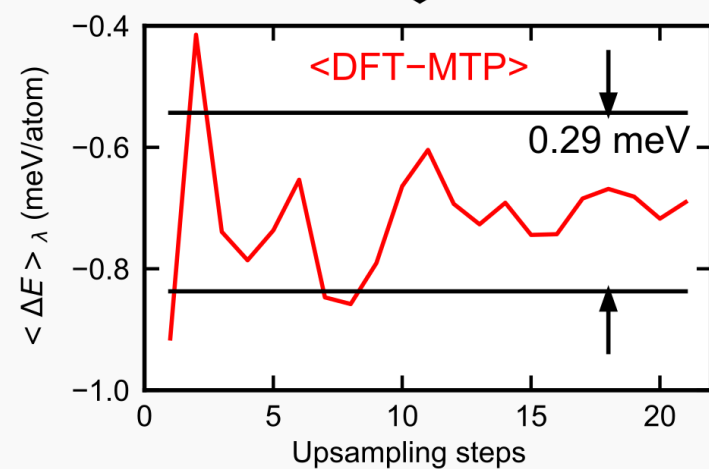
Performed in a large supercell



Upsampling: MTP \rightarrow DFT high

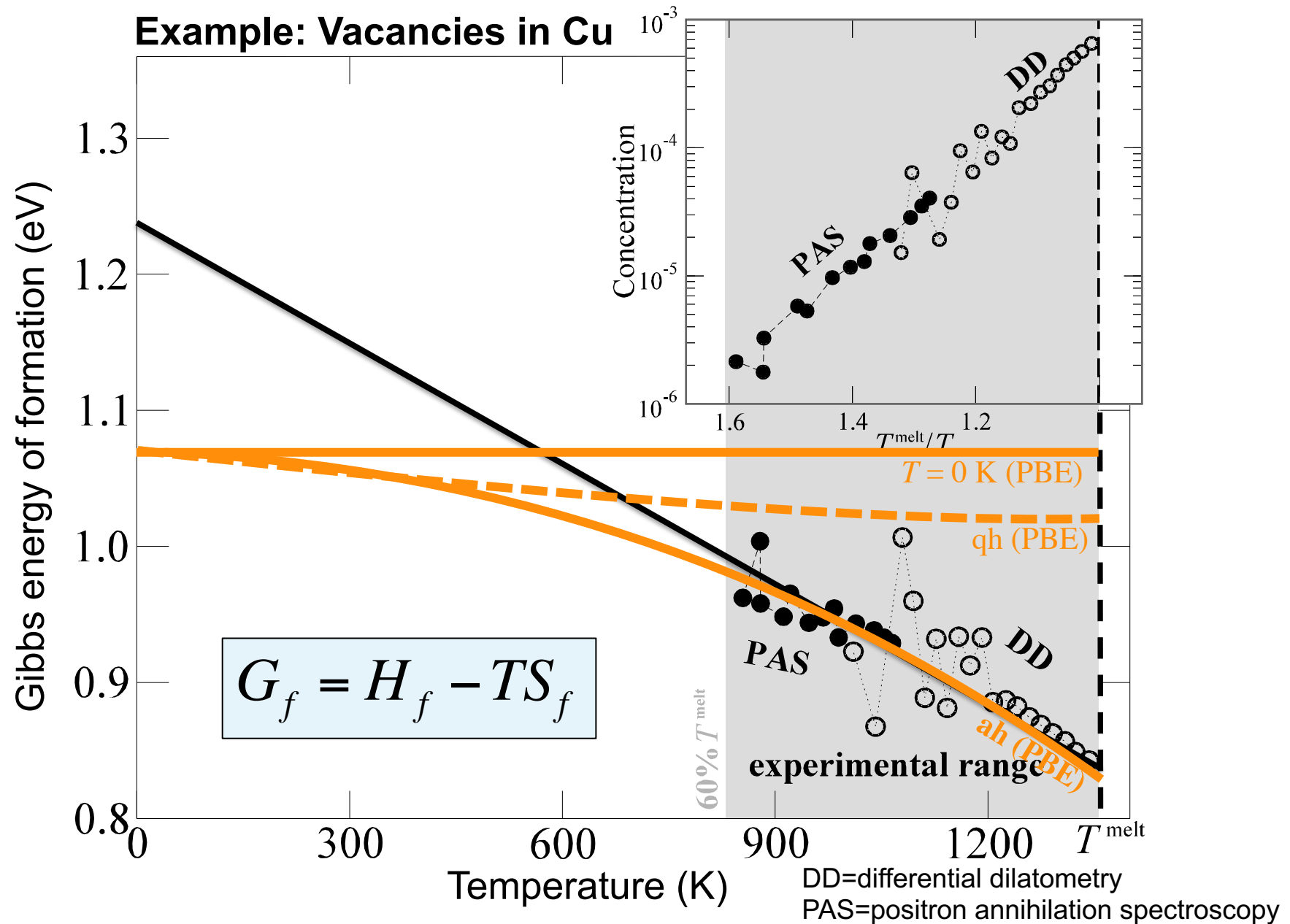


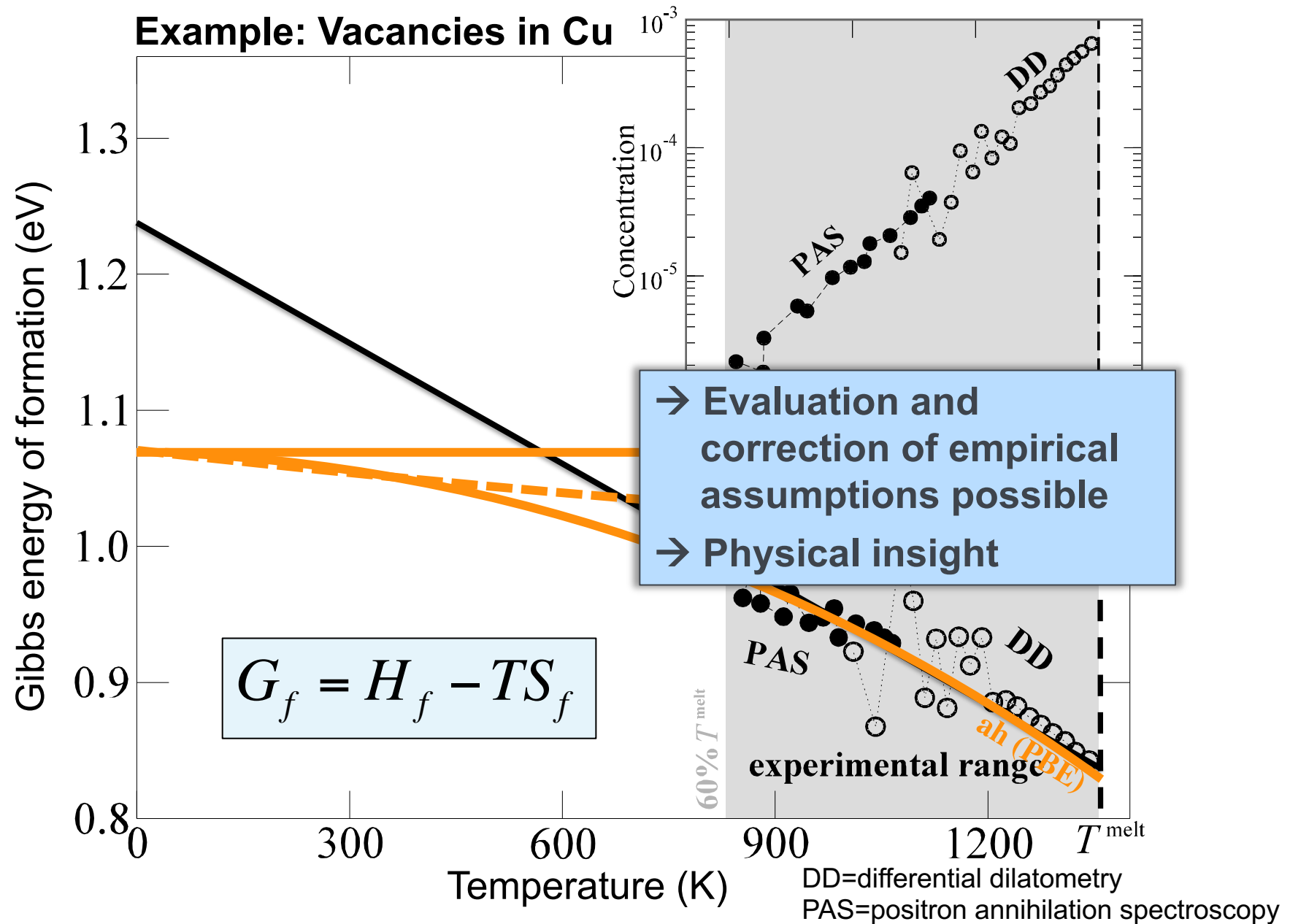
Performed in a medium supercell



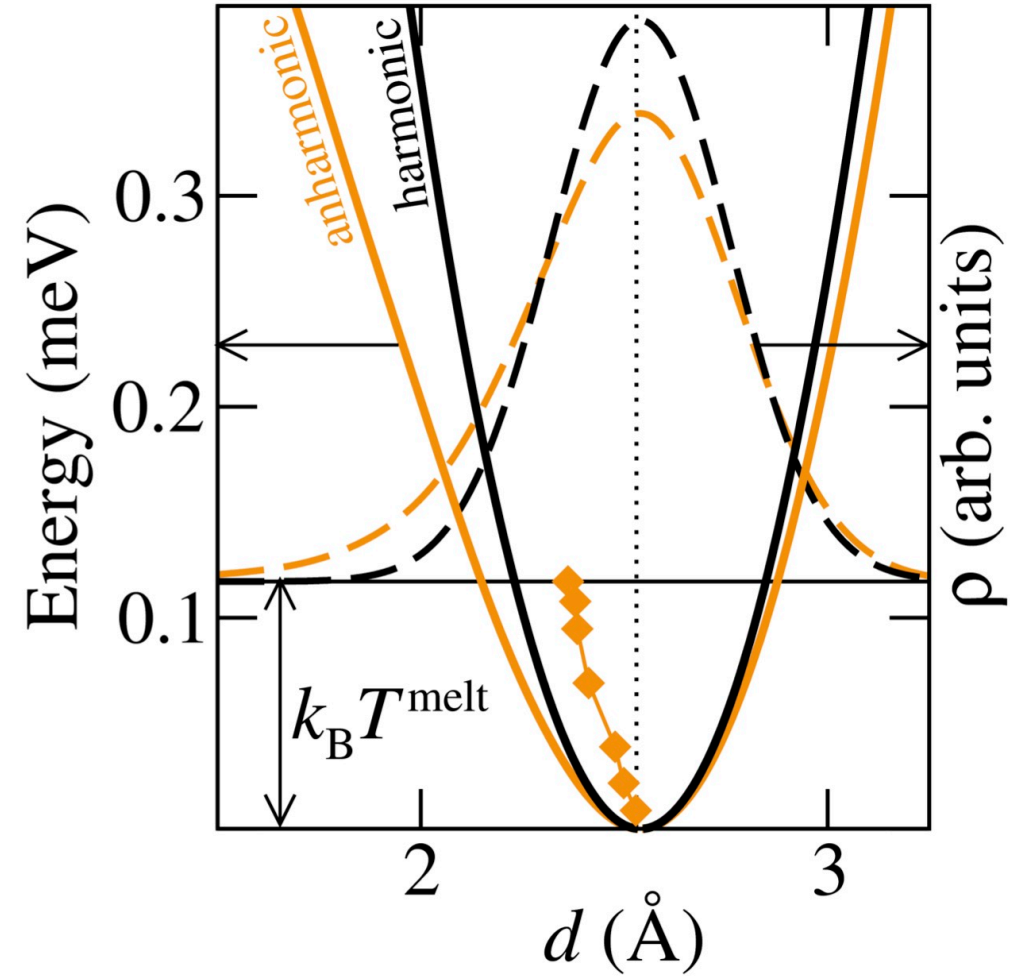
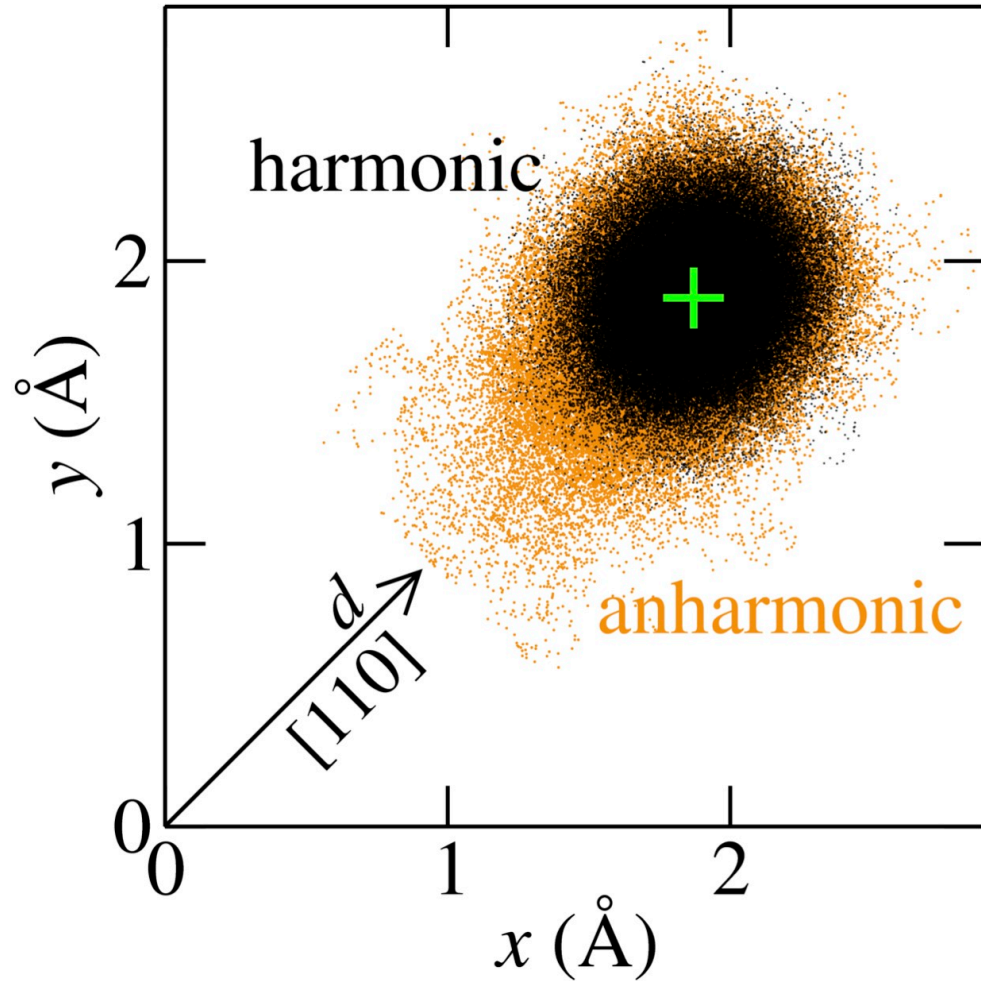
included: anharmonic vibrations, electronic excitations, coupling, magnetism (mean field, empirical)

Defects: Results

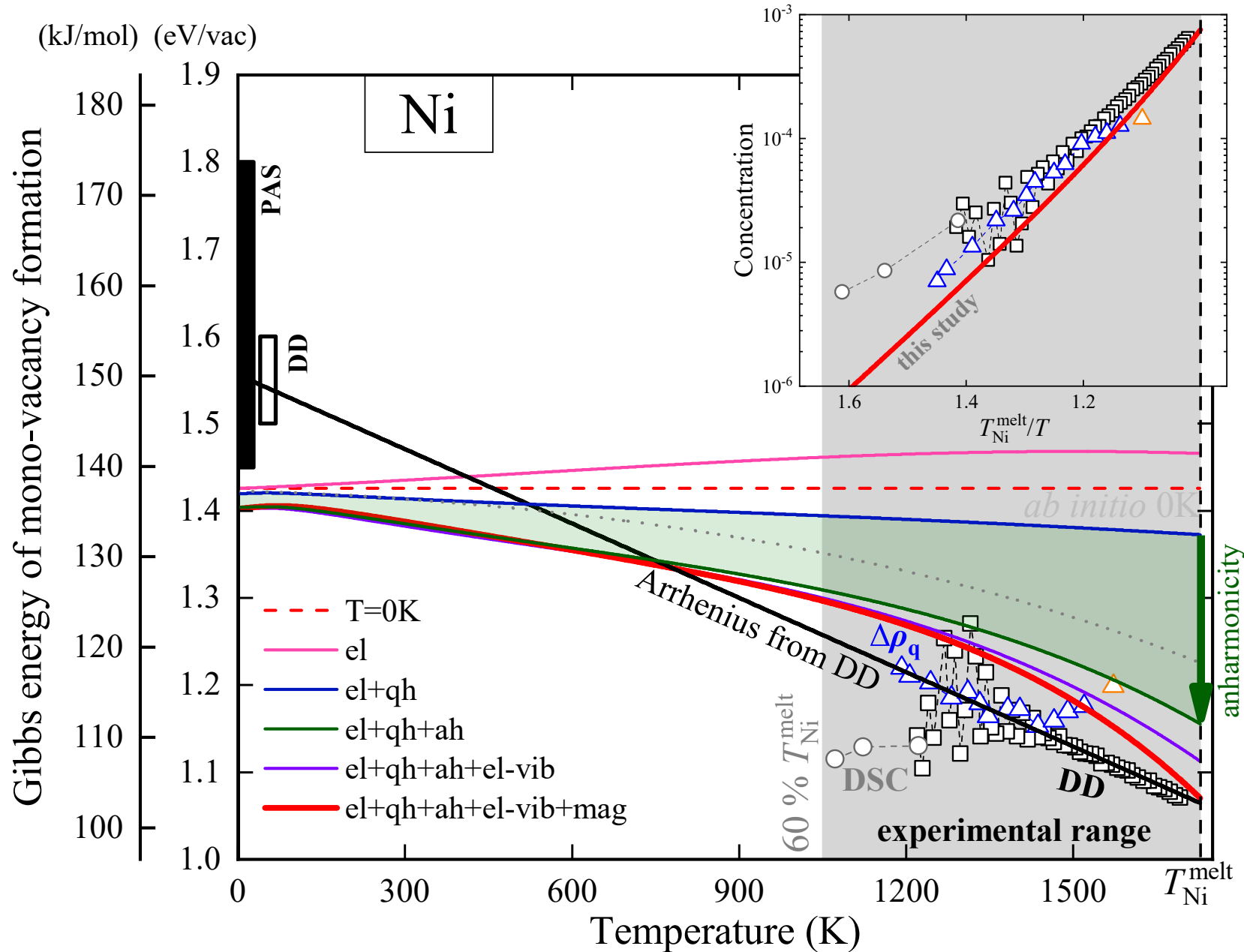


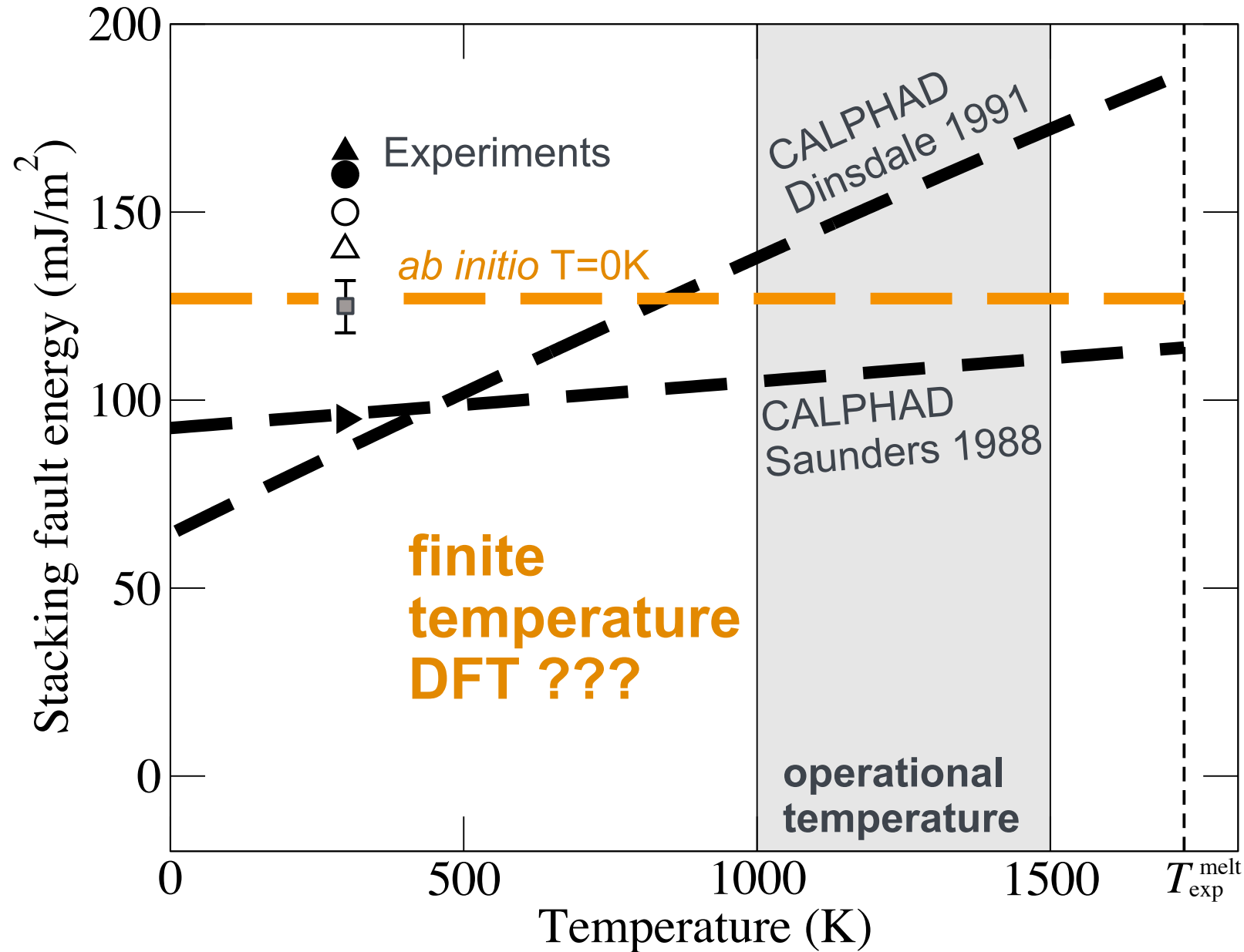


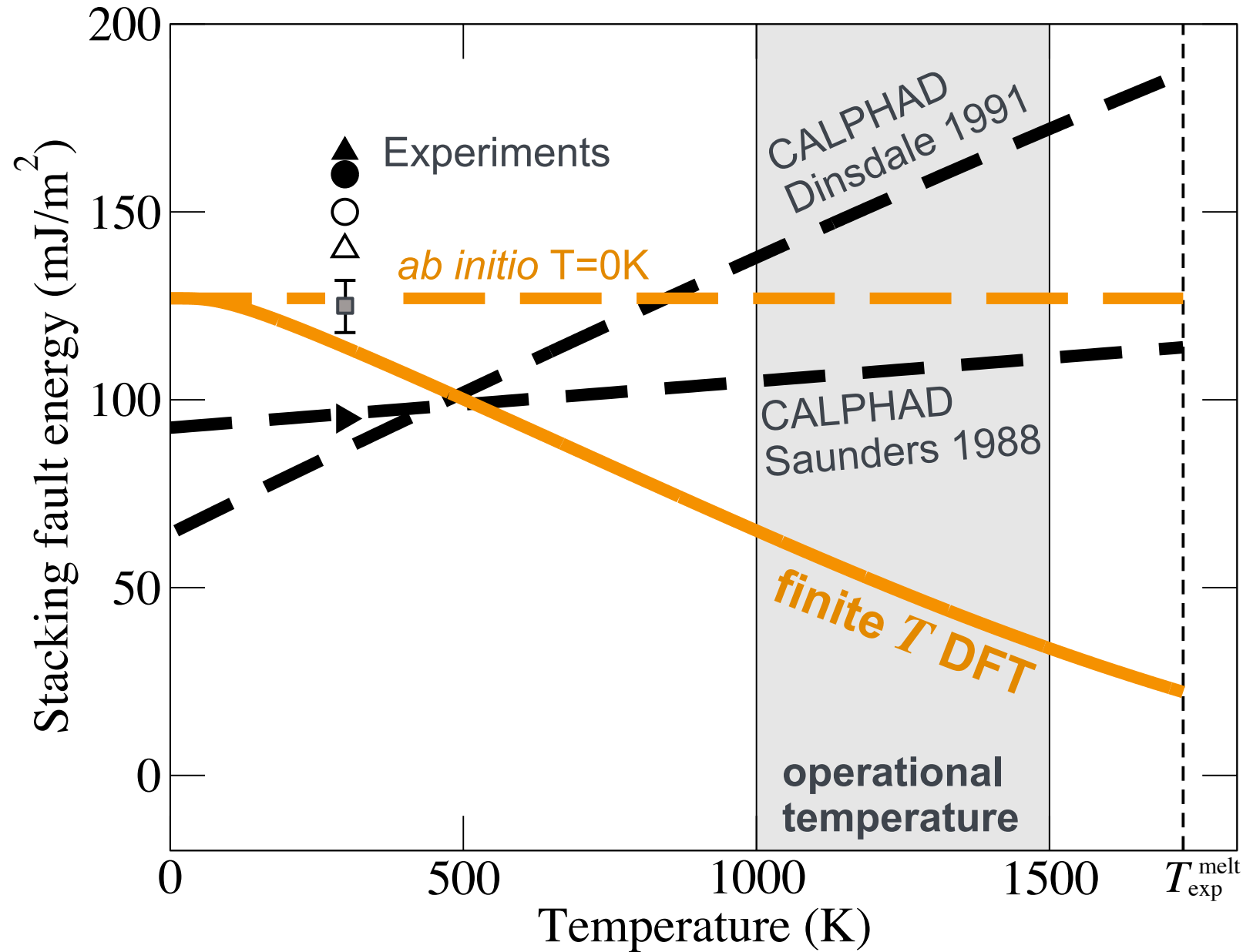
Nearest neighbor distribution at finite temperatures



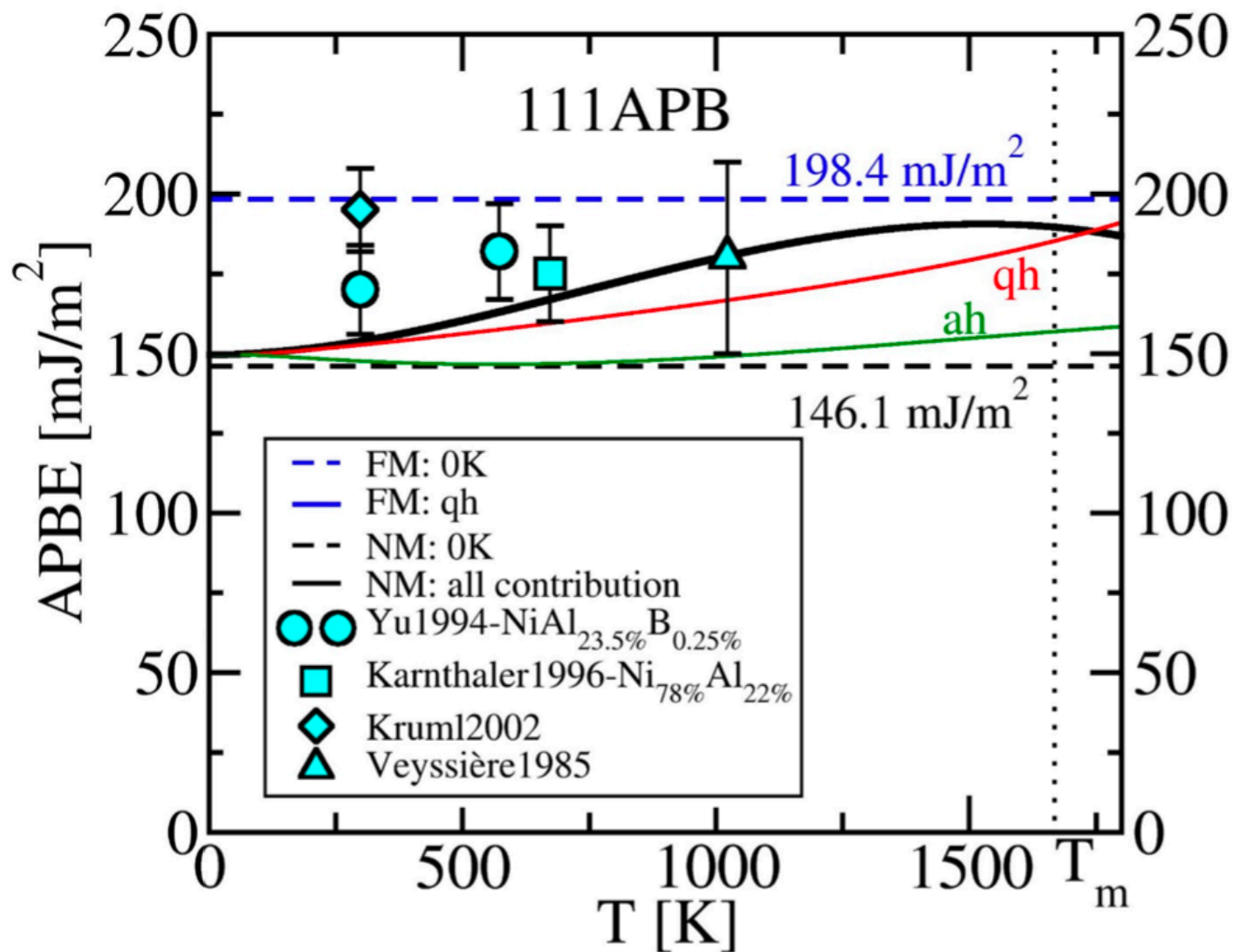
→ Localized interactions around vacancy



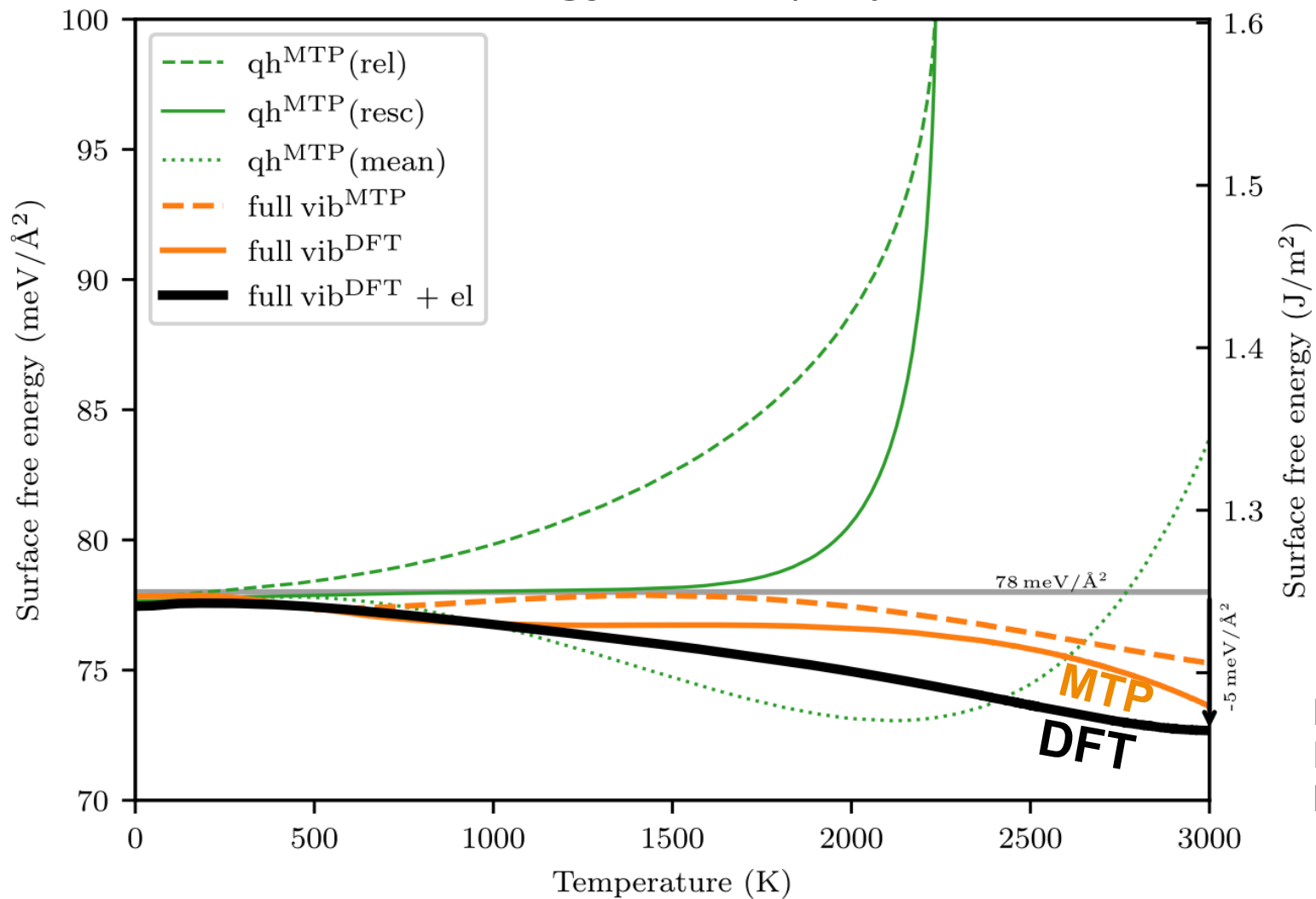




Ni₃Al
APB important
for Ni-superalloys
(to be published)



Surface free energy of TiN (very important for coatings)

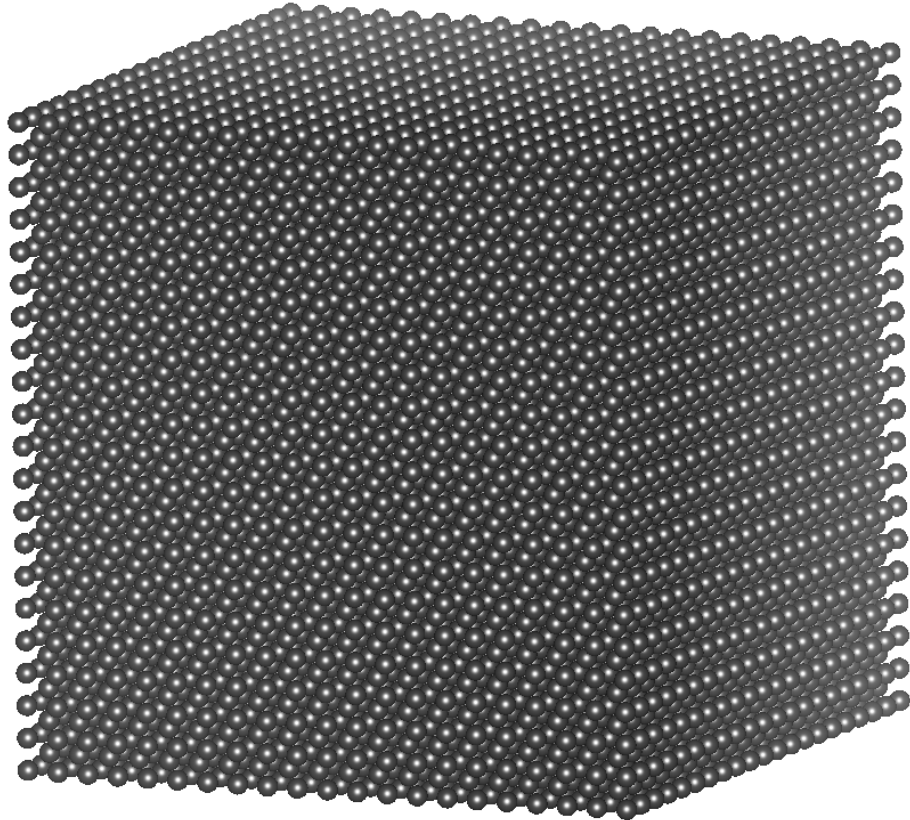


Reference:
Forslund, Grabowski et al.,
PHYSICAL REVIEW B 103,
195428 (2021)

→ high **accuracy** surface properties possible

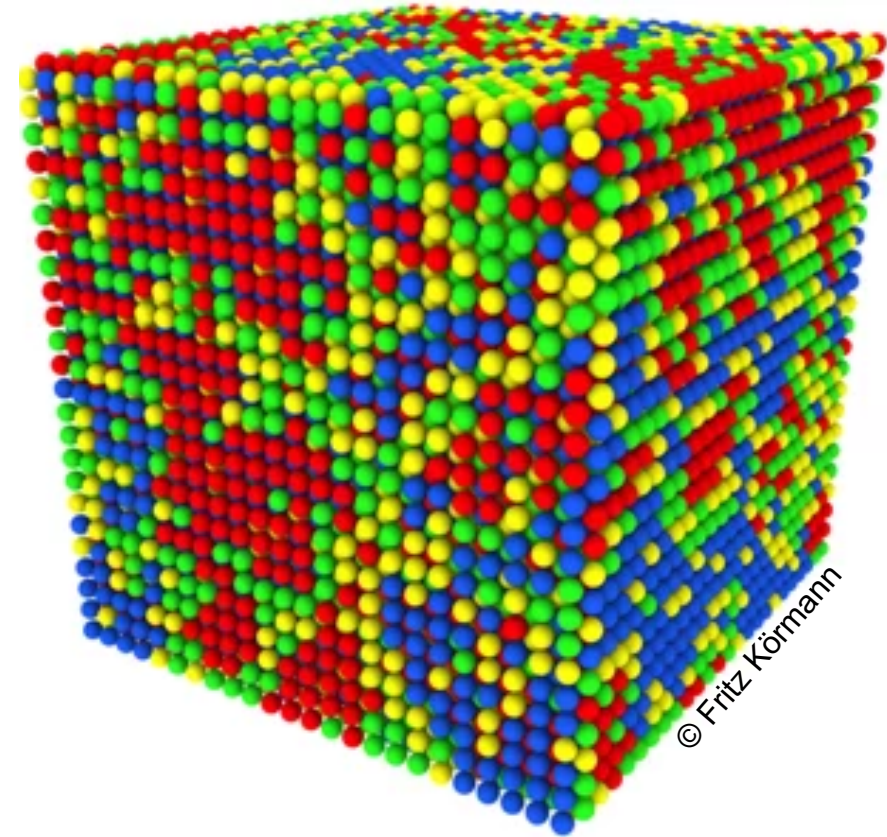
Alloys

Unary system (single element)



→ single defect formation energy

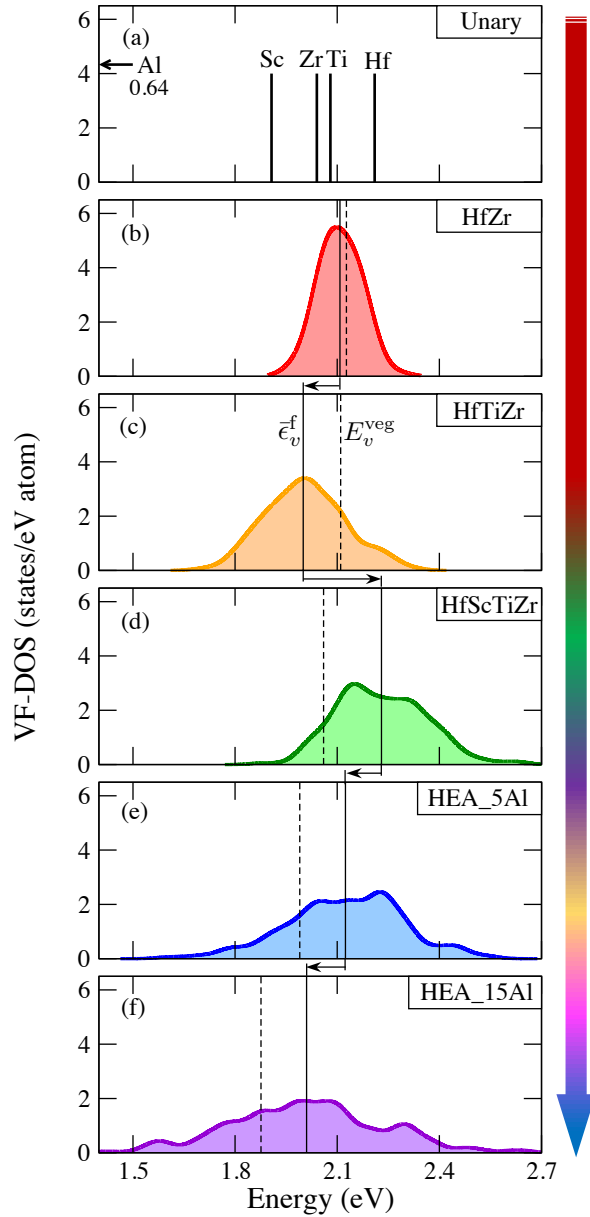
Chemically complex system
High entropy alloy VNbMoTaW



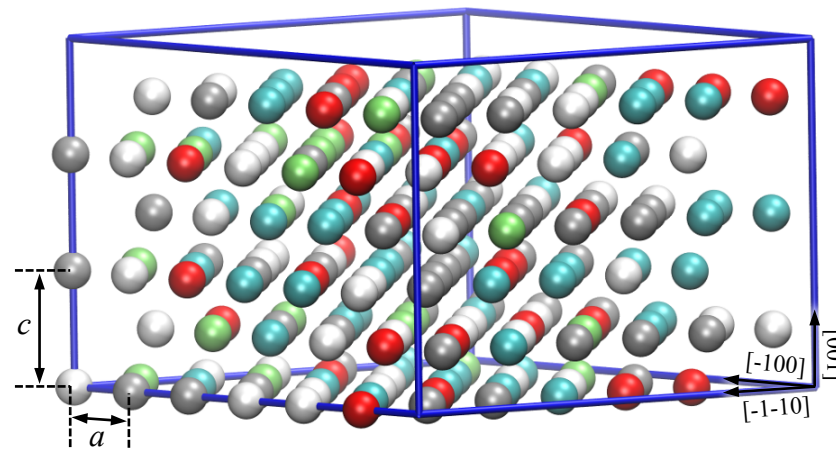
© Fritz Kömann

→ **MANY** defect formation energies (distribution)
→ highly challenging already at 0K

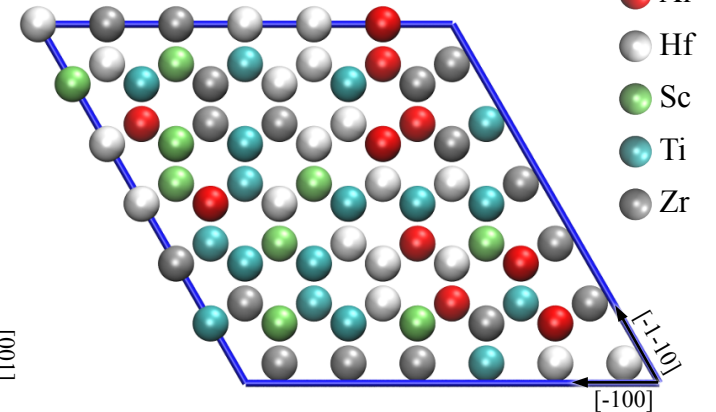
Ref: Xi Zhang, B. Grabowski (under review)



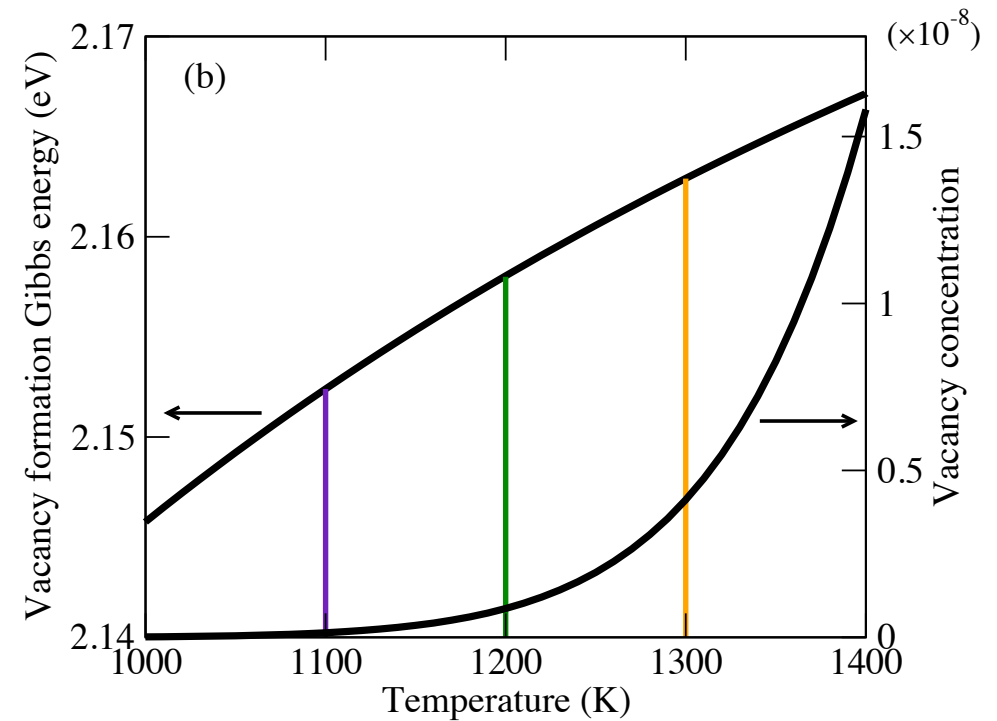
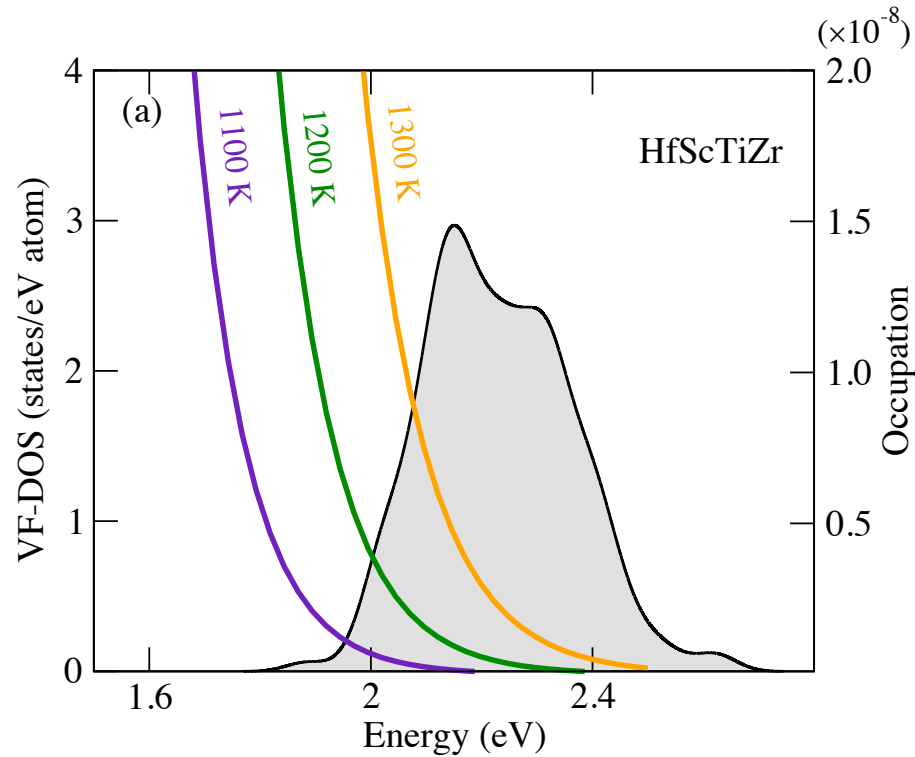
(a) fully disordered hcp A3, side view



(b) fully disordered hcp A3, bottom view

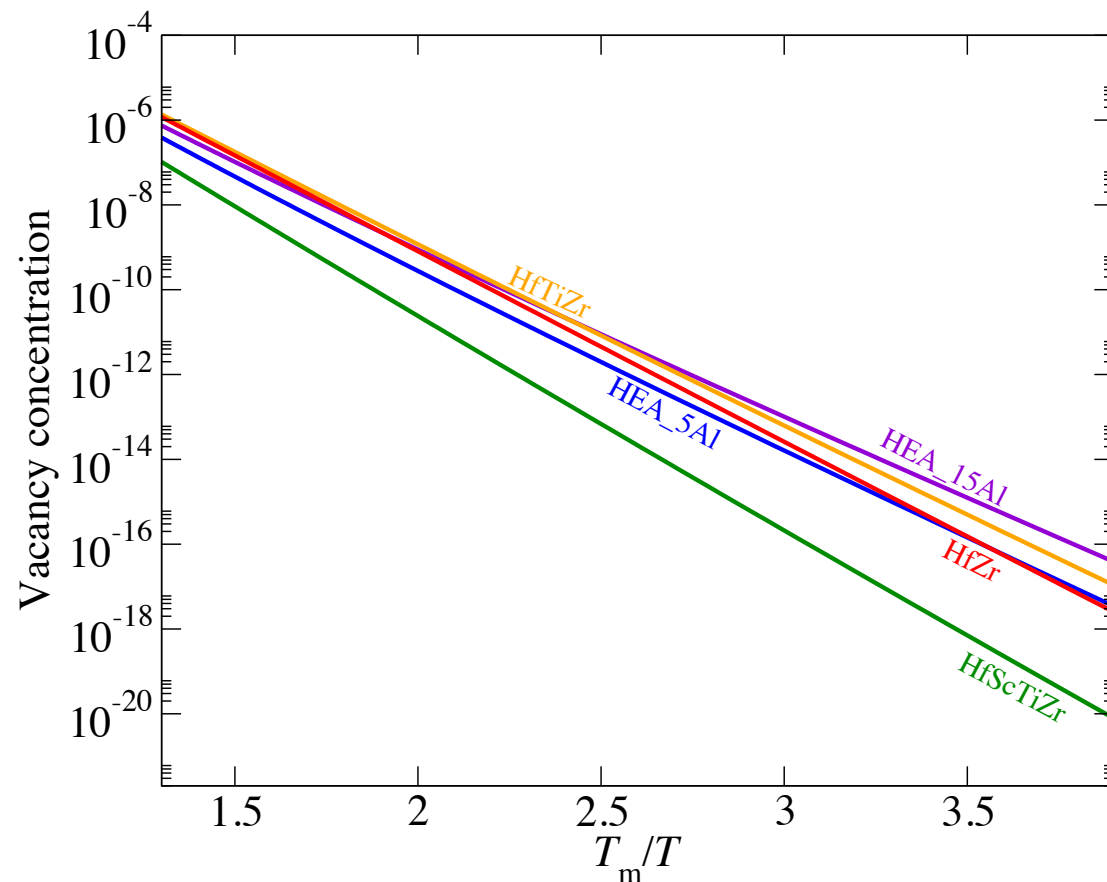
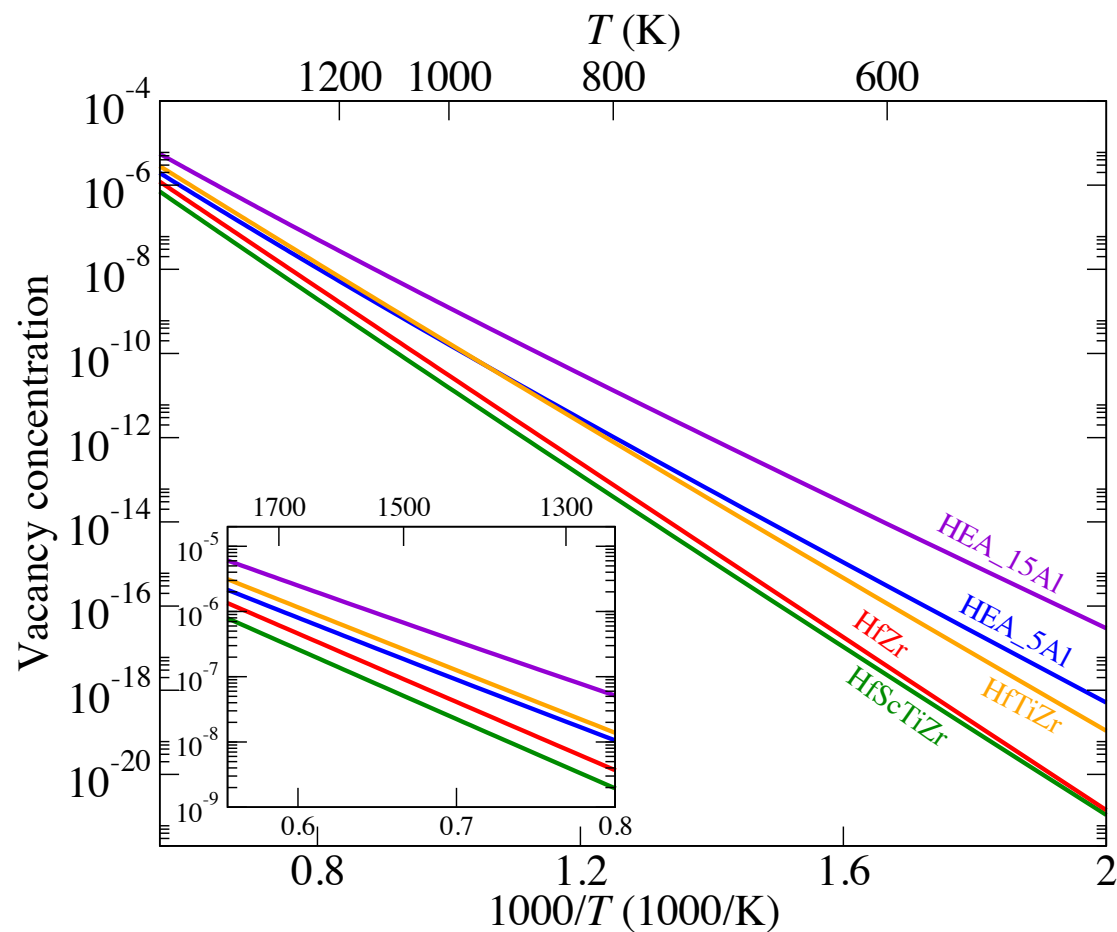


$$x_v^{\text{tot}}(T) = \sum_i x_{v_i} = \sum_i \tilde{D}(\epsilon_{v_i}^f) \exp(-\beta \epsilon_{v_i}^f) \quad \text{occupation function}$$

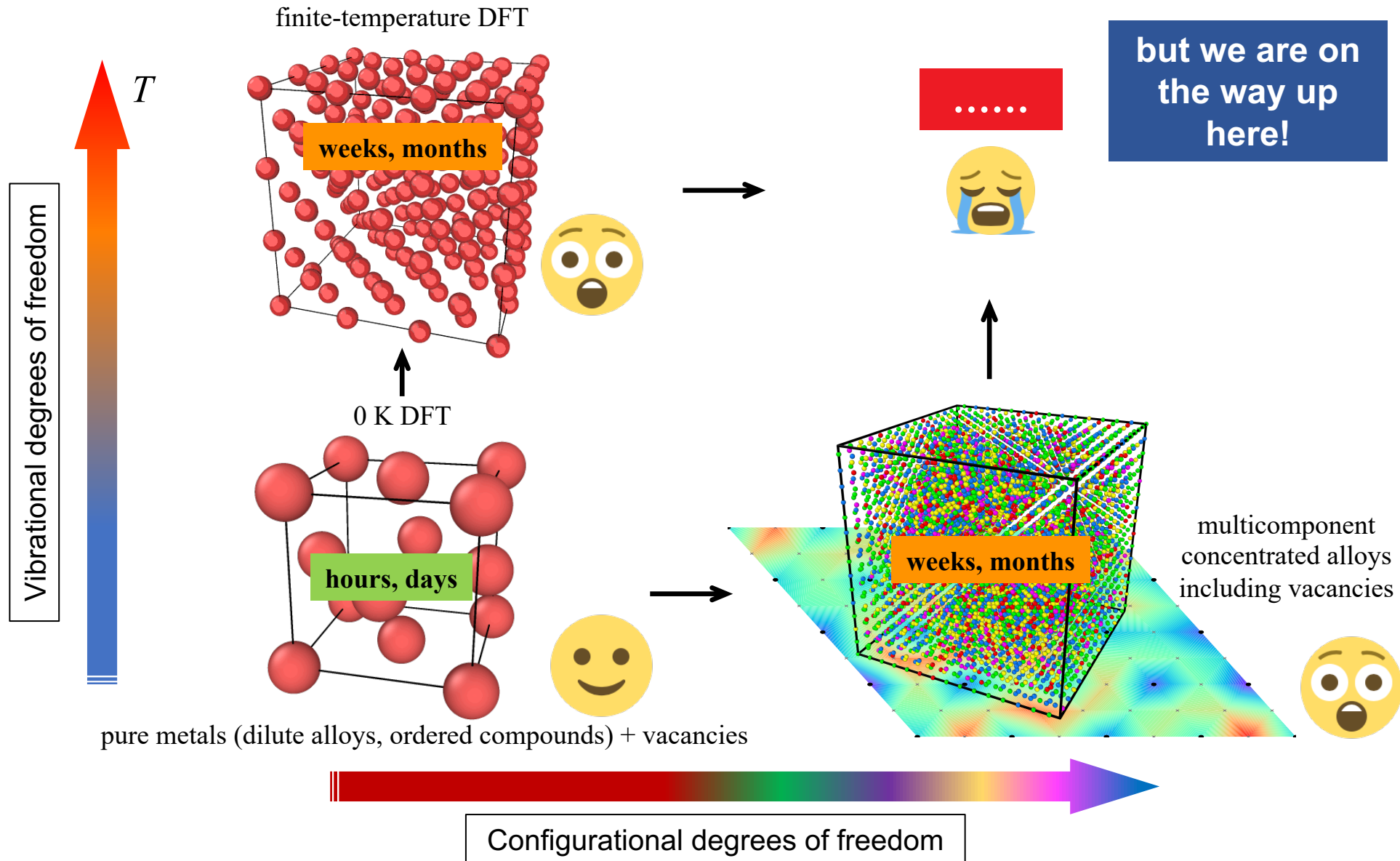


$$x_v^{\text{tot}}(T) = \sum_i x_{v_i} = \sum_i \tilde{D}(\epsilon_{v_i}^f) \exp(-\beta \epsilon_{v_i}^f)$$

non-Arrhenius
dependence



Ref: Xi Zhang, B. Grabowski (under review)



[1] Xi Zhang et al., *Comp. Mater. Sci.* 148 (2018) 249–259

- key: free energy surface, for bulk & defects
- various excitations (qh, ah, el, mag, coupling)
- explicit anharmonic vibrations crucial
- efficient methodology: direct upsampling
- utilizes machine learning potentials
- generally good agreement with experiment
- predictions possible, understanding possible

